# Monotonicity, symmetry and antisymmetry of solutions of semilinear elliptic equations 

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Dedicated to I.M. Gelfand on his seventy-fifth birthday


#### Abstract

This paper investigates certain properties of solutions of equations of the form $$
\Delta u+f(x, u, \nabla u)=0
$$ in a bounded domain $\Omega$ which is convex in the $x_{1}$ direction. Under various conditions on $f$ and on a solution $u$ it is shown that $u$ is increasing in $x_{1}$ in the «left half $\nu$ of $\Omega$, or in all of $\Omega$. Symmetry (in $x_{1}$ ) of some solutions is proved. Also antisymmetry results are obtained. The paper may be considered as an extension of [4].


Key Words: Elliptic partial differential equations. Parabolic equations, maximum principle, monotonicity, symmetry, antisymmetry, uniqueness.

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## 1. INTRODUCTION

In [4] properties of monotonicity and symmetry were established for positive solutions $u$, vanishing on the boundary, of elliptic equations, using the maximum principle and the method of moving planes. The method is due to A.D. Alexandroff and was then used by J. Serrin [8].

A typical result in [4] is the following:
THEOREM A. Let $\Omega$ be a bounded domain in $R^{n}$ with $C^{2}$ boundary which is convex in the $x_{1}$ direction and symmetric about $x_{1}=0$. Let $u$ be a positive solution of

$$
\begin{aligned}
& \Delta u+f(u)=0 \text { in } \Omega \\
& u=0 \text { on } \partial \Omega
\end{aligned}
$$

where $f \in C^{1}$. Then $u_{x_{1}}=u_{1}>0$ for $x_{1}<0$ and $u$ is symmetric in $x_{1}$.
In case $\Omega$ is a ball it then follows that $u$ is radially symmetric and $u_{r}<0$.
The symmetry is a by product of monotonicity in the $x_{1}$ direction-proved by the method of moving planes as follows. For $\lambda$ real let $T_{\lambda}$ be the plane $x_{1}=\lambda$, and

$$
\Sigma(\lambda)=\left\{x \in \Omega ; x_{1}<\lambda\right\}
$$

Let $x^{\lambda}$ denote the reflection of $x$ in $T_{\lambda}$, i.e. $x^{\lambda}$ has the same coordinates as $x$ except for $x_{1}^{\lambda}=2 \lambda-x_{1}$. The monotonicity result which is proved in [4] is

$$
w(x)=u\left(x^{\lambda}\right)-u(x)>0 \text { for } x \in\{\Sigma(\lambda), \lambda<0\}
$$

In the limit we have $u\left(x^{0}\right) \geqslant u(x)$ for $x \in \Sigma(0)$. Since $\Omega$ and the equation are symmetric in $x_{1}$, the opposite inequality follows, and hence symmetry.

Symmetry properties of solutions in all of $R^{n}$ were also proved in [4] and [5].

The authors of this paper extended some of the monotonicity and symmetry results to semilinear equations in infinite cylinders. In the process of doing that they discovered some generalisations of results of [4] in bounded domains. These are the subject of this paper. So this should be regarded as a continuation of [4]. We no longer assume that $u>0$ and $u=0$ on $\partial \Omega$. For convenience we confine ourselves to equations of the form

$$
\begin{equation*}
\Delta u+f(x, u, \nabla u)=0 \tag{1.1}
\end{equation*}
$$

It is always assumed that $f(u, x, p)$ is continuous in all variables and Lipschitz continuous in ( $u, p$ ). The arguments apply to some more general elliptic equations.

In fact after the work was completed Mr. Li, Cong Ming discovered that they could be adapted to handle fully nonlinear second order elliptic equations. In particular he can improve one of the main results, Theorem 2.1' of [4]: he shows that the result holds even if the hypothesis (b) is dropped.

Recently Caffarelli, Gidas and Spruck [11] have used the method of moving planes, rather a measure theoretic variation, to obtain very strong results on behaviour of solutions, of certain elliptic equations, near isolated singular points.

We will now describe some of our results (in $R^{n}$ we represent $x=\left(x_{1}, y\right), y=$ $=\left(x_{2}, \ldots, x_{n}\right)$. In section 2 and 3 we assume that $f$ satisfies for $x_{1}<x_{1}^{\prime}$, $x_{1}+x_{1}^{\prime}<0$,

$$
\begin{equation*}
\hat{f}\left(x_{1}, y, u, p_{1}, \ldots, p_{n}\right) \leqslant f\left(x_{1}^{\prime}, y, u,-p_{1}, p_{2}, \ldots, p_{n}\right) \text { if } p_{1} \geqslant 0 \tag{1.2}
\end{equation*}
$$

Here is a symmetry result which will be proved in section 2 . We consider a finite cylinder

$$
\begin{equation*}
\Omega=S_{a}=\left\{\left(x_{1}, y\right) \in R^{n} ;\left|x_{1}\right|<a, y \in \omega\right\} \tag{1.3}
\end{equation*}
$$

where $\omega$ is a bounded domain in $R^{n-1}$ with smooth boundary.
THEOREM 1.1. Let $u$ be a $C^{2}$ solution of (1.1) in $\bar{\Omega}$ and assume that the boundary values of $u$ are symmetric in $x_{1}$ and satisfy

$$
\begin{equation*}
u_{1}\left(x_{1}, y\right) \geqslant 0 \text { for }-a<x_{1}<0, \quad y \in \partial \omega \tag{1.4}
\end{equation*}
$$

## Assume also that

$$
\left\{\begin{array}{l}
u(-a, y) \leqslant u\left(x_{1}, y\right) \text { for }-a<x_{1}<a, y \in \omega  \tag{1.5}\\
\text { and } \forall x_{1} \text { in }(-a, a), \exists y \in \omega \text { such that strict inequality holds. }
\end{array}\right.
$$

Assume that $f$ satisfies (1.2) and that it is symmetric in $\left(x_{1}, p_{1}\right)$. Then $u$ is a symmetric function of $x_{1}$ and $u_{1}>0$ if $x_{1}<0$.

The theorem will be derived from a monotonicity result, Theorem 2.1, for solutions $u$ of (1.1) in $S_{a}$ satisfying (1.5). Theorem 2.2 contains a similar result under Neumann boundary conditions on the curved part of $\partial S_{a}$.

Here is an example in which the theorem applies. In the cylinder $\Omega$ let $z$ be a solution of

$$
\begin{aligned}
& \Delta z+g(z)=0 \\
& z=\phi \text { on } \partial \Omega
\end{aligned}
$$

where $\phi \geqslant 0, \phi$ is symmetric in $x_{1}, \phi_{1} \geqslant 0$ for $x_{1}<0, y \in \partial \omega$, and $\phi=0$ at the ends: $x_{1}= \pm a$. We assume $g$ is a $C^{1}$ nonnegative function. For $\sigma(y)$ a smooth function
of $y$, the function $u=z-\sigma(y)$, satisfies (1.1) with $f=g(u+\sigma(y))+\Delta_{y} \sigma$.
In section 3 we extend the $x_{1}$-monotonicity results to general domains $\Omega$ with smooth boundary which are convex in the $x_{1}$ direction. Again we replace the conditions of [4], that $u>0$ in $\Omega, u=0$ in $\partial \Omega$, by the condition that on any segment in $\Omega$ parallel to the $x_{1}$ axis, $u$ is greater than its value at the left end point lying on $\partial \Omega$. In addition, assuming $u$ is less than its value at the right end point in $\partial \Omega$ we prove monotonicity (in $x_{1}$ ) in the full domain $\Omega$.

Here is a special case of Theorem 3.1':
THEOREM 1.1': Let $\Omega$ be a bounded domain, with smooth boundary, which is convex in the $x_{1}$ direction and symmetric in $x_{1}$. In $\bar{\Omega}$ let u be a $C^{3}$ solution of

$$
\begin{aligned}
& \Delta u+f(u)=0 \\
& u=\phi \quad \text { on } \quad \partial \Omega
\end{aligned}
$$

where $f$ is Lipschitz continuous Assume that $\phi$ is symmetric in $x_{1}$ and that for $x_{1} \leqslant 0$, wherever the normal $\nu$ to $\partial \Omega$ has $\nu_{1}=0$, there $\phi_{1} \leqslant 0$. Assume also that if $\left(x_{1}, y\right) \in \Omega,\left(x_{1}^{\prime}, y\right) \in \partial \Omega$ then

$$
u\left(x_{1}, y\right)>\phi\left(x_{1}^{\prime}, y\right)
$$

Then $u$ is symmetric in $x_{1}$ and $u_{1}>0$ for $x_{1}<0$.
L. Caffarelli indicated to us a different and very simple argument for proving $x_{1}$-monotonicity under the condition (in place of (1.2)): $f$ is nondecreasing in $x_{1}$. See Theorem 3.5 and its corollary. In section 4 we modify his argument to prove monotonicity under the weaker condition

$$
\begin{equation*}
f(x, u, p) \text { is nondecreasing in } x_{1} \text { for } p_{1} \geqslant 0 \tag{1.6}
\end{equation*}
$$

Here is a simplified form of Theorem 4.1 in the cylinder $\Omega=S_{a}$.
THEOREM 1.2. Let $u \in C^{2}(\bar{\Omega})$ be a solution of (1.1) with $f$ satisfying (1.6) and Lipschitz continuous in $\left(x_{1}, u, p\right)$. Assume that $u_{1} \geqslant 0$ on $\partial \Omega$ and that

$$
\begin{equation*}
u(-a, y)<u\left(x_{1}, y\right)<u(a, y) \quad \forall y \in \omega,-a<x_{1}<a \tag{1.5}
\end{equation*}
$$

Then $u_{1}>0$ in $\Omega$. Furthermore if $\underline{u}$ is any solution of (1.1) agreeing with $u$ on $\partial \Omega$ and satisfying (1.5)', then $\underline{u}=u$.

We conclude the introduction with some remarks about the maximum principle for solutions in a domain $\Omega$ of an elliptic inequality (we use summation convention)

$$
\begin{equation*}
a_{i j}(x) u_{x_{i} x_{j}}+b_{i}(x) u_{x_{i}}+c(x) u \geqslant 0 . \tag{1.7}
\end{equation*}
$$

Uniform ellipticity is assumed:

$$
\begin{equation*}
a_{i j} \xi_{i} \xi_{j} \geqslant c_{0}|\xi|^{2}, c_{0}>0 \tag{1.8}
\end{equation*}
$$

We will constantly refer to results from [4]: the Maximum Principle and its Corollary as well as Lemma $H$ on pages 212-3 of [4], and also to the extended Hopf boundary lemma at a corner, Lemma S, in [4]. In addition we will use Lemma A. 1 there. In all of these results, as stated in [4], it is assumed that $u \in C^{2}(\bar{\Omega})$ and that the coefficients in (1.7) are continuous (for Lemma $S$ further regularity of the $\left\{a_{i j}\right\}$ is required).

For various applications it is important to be able to relax these smoothness requirements. In [1] Amick and Fraenkel have extended (and used) these results for equations in divergence form, with merely bounded measurable coefficients. For equations (1.7) in nondivergence form we point out that it is enough to suppose that $u \in W^{2, p}, p>n$, and that the coefficients in (1.7) belong to $L^{\infty}$. Namely, the Maximum Principle and its Corollary, and Lemma H of [4] hold for such solutions. The proofs of these assertions proceed as in the classical cases, with the aid of the usual barrier functions, but using the Bony maximum principle [3]. Recall its statement (see P.L. Lions [6] for an extension):

Bony Max. Princ.: Let $u \in W^{2, p} p>n$, in a neighbourhood of the origin in $R^{n}$, and have a local maximum at the origin. Then

$$
\lim _{x \rightarrow 0} \inf \text { ess } \Sigma a_{i j}(x) u_{i j}(x) \leqslant 0
$$

where $\left\{a_{i j}\right\}$ is a nonnegative matrix belonging to $L_{\text {loc }}^{\infty}(\Omega)$.
Furthermore, Lemma $S$ of [4] holds for $C^{2}(\bar{\Omega})$ solutions of (1.7) in case $a_{i j} \in C^{2}(\bar{\Omega})$, and the coefficients $b_{i}$ and $c$ are in $L^{\infty}$.

New ideas are involved in proving the new results. A new ingredient in our application of the method of moving planes is the use of parabolic inequalities and the corresponding maximum principle. This came as a surprise to us. In particular we will need an analogue of the corollary on page 213 of [4]. We state it only in the simplest form. In $R^{n+1}$ with coordinates $(x, t), x \in R^{n}$, let $V$ be a bounded domain lying in $t<T$. In $\bar{V}$ we consider a function $w$, with continuous derivatives up to second order in $x$ and first order in $t$, satisfying a (degenerate) parabolic inequality:

$$
\begin{equation*}
\left(L-\beta \partial_{t}\right) w=a_{i j}(x, t) w_{x_{i} x_{j}}+b_{i}(x, t) w_{x_{i}}+c(x, t) w-\beta(x, t) w_{t} \geqslant 0 . \tag{1.9}
\end{equation*}
$$

Here the $a_{i j}$ are continuous and satisfy (1.8) and the other coefficients are all in $L^{\infty}$; in addition $\beta \geqslant 0$.

PROPOSITION 1.1. Assume that at every point of

$$
J=\overline{\partial V \cap\{t<T\}}
$$

we have $w \leqslant 0$. If $V$ lies in a sufficiently narrow region $0<a+x_{1}<\epsilon$ then $w \leqslant 0$ in $V$.

Here $\epsilon$ depends in the constant $c_{0}$ in (1.8) and on the $L^{\infty}$ norms of the coefficients $b_{i}, c, \beta$.

The proposition will be proved in the Appendix together with analogues of Lemmas H and S which are stated in section 4 . See also Lemma A.1.

We wish to express our thanks to $\mathbf{L}$. Caffarelli for several useful remarks and to Li , Cong Ming for his remarks and simplifications of some of the proofs.

## 2. MONOTONICITY

In this section we will present generalizations of one of the basic results, Theorem 2.1 of [4] (1). The reader should note however that, here, the moving planes will move in the direction of increasing $x_{1}$ (in [4] it was the other way). So our conditions will look a bit different from those in [4]. As we have remarked, we do not assume $u=0$ on $\partial \Omega$ and $u>0$ in $\Omega$. Rather we assume, essentially, that on any open interval in $\Omega$ parallel to the $x_{1}$ axis $u$, is greater than its value at the left end point (i.e. smallest $x_{1}$ ) on $\partial \Omega$.

In this section we consider the simplest geometry, the finite cylinder $\Omega=S_{a}$ of $\S 1$ :

$$
\begin{equation*}
S_{a}=\left\{\left(x_{1}, y\right) \in R^{n} ;\left|x_{1}\right|<a, y \in \omega\right\} \tag{2.1}
\end{equation*}
$$

In $[-a, a) \times \bar{\omega}$ we consider a $C^{2}$ solution of

$$
\begin{equation*}
\Delta u+f(x, u, \nabla u)=0 \tag{2.2}
\end{equation*}
$$

We will prove monotonicity of $u$ in $x_{1}$ for $x_{1}<0$ assuming (1.2):

$$
\left\{\begin{array}{l}
\text { for } x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0, \text { and } p_{1} \geqslant 0,  \tag{2.3}\\
f\left(x_{1}, y, u, p_{1}, p_{2}, \ldots, p_{n}\right) \leqslant f\left(x_{1}^{\prime}, y, u,-p_{1}, p_{2}, \ldots, p_{n}\right)
\end{array}\right.
$$

Observe that if $p_{1}<0$ then, by Lipschitz continuity, $\exists C \geqslant 0$ such that for $x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0$,

[^0]\[

$$
\begin{aligned}
& f\left(x_{1}^{\prime}, y, u,-p_{1}, p_{2}, \ldots, p_{n}\right)-f\left(x_{1}, y, u, p_{1}, p_{2}, \ldots, p_{n}\right) \\
& \geqslant f\left(x_{1}^{\prime}, y, u, 0, p_{2}, \ldots, p_{n}\right)-f\left(x_{1}, y, u, 0, p_{2}, \ldots, p_{n}\right)+C p_{1} \\
& \geqslant C p_{1} \text { by (2.3). }
\end{aligned}
$$
\]

It follows that there is an $L^{\infty}$ function $\beta \geqslant 0$, such that for $x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0$, and all $p_{1}$,

$$
\begin{equation*}
f\left(x_{1}^{\prime}, y, u,-p_{1}, p_{2}, \ldots, p_{n}\right)-f\left(x_{1}, y, u, p_{1}, p_{2}, \ldots, p_{n}\right) \geqslant \beta p_{1} \tag{2.4}
\end{equation*}
$$

Here $\beta$ depends on $x_{1}, x_{1}^{\prime}, p$, etc.
THEOREM 2.1. With $u$ and $f$ as above, assume

$$
\begin{align*}
& u(-a, y) \leqslant u\left(x_{1}, y\right) \text { for }-a<x_{1}<a, y \in \omega  \tag{2.5}\\
& \text { and, } \forall x_{1} \text { in }(-a, a), \exists y \in \omega \text { such that strict inequality holds. }
\end{align*}
$$

In addition we assume that for $y \in \partial \omega,-a<x_{1}<x_{1}^{\prime}<a$

$$
\begin{equation*}
u\left(x_{1}, y\right) \leqslant u\left(x_{1}^{\prime}, y\right) \text { provided } x_{1}+x_{1}^{\prime}<0 \tag{2.6}
\end{equation*}
$$

Then, for $-a<x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0, y \in \omega$ we have

$$
\begin{equation*}
u_{1}\left(x_{1}, y\right)>0, u\left(x_{1}, y\right)<u\left(x_{1}^{\prime}, y\right) \tag{2.7}
\end{equation*}
$$

Furthermore if $u_{1}(0, y)=0$ for some $y \in \omega$ then $u$ is a symmetric function of $x_{1}$.

REMARK 2.1. We have stated the theorem in case $n>1$, but of course it holds also for $n=1$ - simply ignore the boundary conditions on $\partial \omega$. Here is a simple example in case $n=1$ showing that the condition (2.3) on $f$ cannot be dropped: On the interval $\Omega=(-a, a)$, a large, the conclusion of Theorem 2.1 does not hold for

$$
u=(x+a) e^{-x}
$$

## It satisfies

$$
\ddot{u}+2 \dot{u}+u=0
$$

and condition (2.3) does not hold. Furthermore if $u$ is any $C^{3}$ function on $[-a, a]$ with $u(-a)<u(x)<u(a)$ which is not monotone, it still satisfies a differential equation $\ddot{u}+f(x)=0$ where $f$ is in fact $-\ddot{u}(x)$.

Observe also (see (2.5)) that on the interval $|x|<a=3 \pi / 2$ the function

$$
\begin{equation*}
u=1-\sin x \text { satisfies } \tag{2.8}
\end{equation*}
$$

$$
\ddot{u}+u-1=0
$$

and satisfies $u(-a)<u(x)<u(a)$ for $|x|<a$ except at $x= \pm \pi / 2$, and (2.7) does not hold for $u$. For $n=2$ we may take $a=3 \pi / 2, \omega=(0, \pi)$ and $u=1$ -$-\sin x \sin y$, which satisfies $\Delta u+2 u-2=0$. Condition (2.6) holds as does (2.5) except at $x=\pi / 2$, and (2.7) does not hold.

The proof of the Theorem is similar to those of Theorems 2.1 and 3.1 in [4] but differs in some essential details. In particular we do not have the analogue of Lemma 2.1 of [4].

We will use similar notation: For any $\lambda$ in $-a<\lambda \leqslant 0$ let $\Sigma(\lambda)$ denote the finite cylinder

$$
\Sigma(\lambda)=\left\{\left(x_{1}, y\right) \in \Omega ;-a<x_{1}<\lambda\right\} .
$$

$\Sigma^{\prime}(\lambda)$ is the reflection of $\Sigma(\lambda)$ in the plane

$$
T_{\lambda}=\left\{x_{1}=\lambda\right\} ;
$$

the reflection of $x=\left(x_{1}, y\right)$ in the plane $T_{\lambda}$ is the point $x^{\lambda}=\left(2 \lambda-x_{1}, y\right)$.

Proof of Theorem $2.1: \operatorname{In} \Sigma(\lambda)$ consider the functions

$$
\begin{equation*}
v(x)=u\left(x^{\lambda}\right)=u\left(2 \lambda-x_{1}, y\right), \text { and } w(x)=w_{\lambda}(x)=v(x)-u(x) \tag{2.9}
\end{equation*}
$$

We will prove that for every $\lambda$ in $(-a, 0)$,

$$
\begin{equation*}
w(x)>0 \text { in } \Sigma(\lambda) \tag{2.10}
\end{equation*}
$$

These yield (2.7). We use the method of moving planes; it consists of two steps: (I) Initial step: prove (2.10), (2.11) for $0<a+\lambda$ small.
(II) Continuation: prove the inequalities for all $\lambda$ in $(-a, 0)$.
(I) Here we use condition (2.4) to derive a parabolic differential inequality for $w$ in $\Sigma(\lambda)$. There $v$ satisfies

$$
\begin{aligned}
& -\Delta v=f\left(x^{\lambda}, v,-v_{1}, \nabla_{y} v\right) . \\
& \geqslant f\left(x, v, v_{1}, \nabla_{y} v\right)+\beta v_{1}
\end{aligned}
$$

with $\beta \in L^{\infty}, \beta(x, \lambda) \geqslant 0$, by (2.4). Hence $w$ satisfies

$$
\begin{aligned}
& -\Delta w=f\left(x^{\lambda}, v,-v_{1}, \nabla_{y} v\right)-f(x, u, \nabla u) \\
& \geqslant f(x, v, \nabla v)-f(x, u, \nabla u)+\beta v_{1}
\end{aligned}
$$

Since $f(x, u, p)$ is Lipschitz continuous in $(u, p)$ it follows that for suitable bounded functions $b_{j}, c$,

$$
\Delta w+\sum_{1}^{n} b_{j} w_{j}+c w+\beta v_{1} \leqslant 0
$$

But

$$
\frac{\partial}{\partial \lambda} w=2 u_{1}\left(x^{\lambda}, y\right)=-2 v_{1}(x)
$$

Thus we obtain a (degenerate) parabolic inequality for $w$ as a function of $x$ and $\lambda$ :

$$
\begin{equation*}
\Delta w+b_{j} w_{j}+c w-\frac{\beta}{2} \frac{\partial}{\partial \lambda} \quad w \leqslant 0 \tag{2.12}
\end{equation*}
$$

It holds in a region $V$ in $(x, \lambda)$ space:

$$
V=\left\{\left(x_{1}, y, \lambda\right) ;-a<x_{1}<\lambda<\lambda_{0}, y \in \omega\right\}
$$

Except on the top part, $\lambda=\lambda_{0}$, we have $w \geqq 0$ on $\partial V$. Indeed on the boundary where $x_{1}=\lambda$ we have $w=0$, and where $x_{1}=-a$ we have $w(x, \lambda) \geqq 0$ by (2.5). For $0<a+\lambda_{0}$ small, the width of this region in the $x_{1}$ direction is small, namely $a+\lambda_{0}$. We may therefore apply Proposition 1.1 and conclude that

$$
w \geq 0 \text { in } \Sigma(\lambda),
$$

and hence

$$
-2 u_{1}=w_{1} \leqslant 0 \text { on } T_{\lambda} \cap \Omega
$$

this is true for every $\lambda$ close to $-a$. So $u_{1} \geqslant 0$ for $x_{1}$ close to $-a$.
To finish Step (I) we use (2.3) and derive an elliptic inequality for $w$ in $\Sigma(\lambda)$ : As before we have

$$
\begin{aligned}
& -\Delta w=f\left(x^{\lambda}, v,-v_{1}, \nabla_{y} v\right)-f\left(x, u, u_{1}, \nabla_{y} u\right) \\
& =f\left(x^{\lambda}, v,-v_{1}, \nabla_{y} v\right)-f\left(x^{\lambda}, u,-u_{1}, \nabla_{y} u\right) \\
& +f\left(x^{\lambda}, u,-u_{1}, \nabla_{y} u\right)-f\left(x, u, u_{1}, \nabla_{y} u\right) \\
& \geqslant f\left(x^{\lambda}, v,-v_{1}, \nabla_{y} v\right)-f\left(x^{\lambda}, u,-u_{1}, \nabla_{y} u\right)
\end{aligned}
$$

by (2.3) (we have shown that $u_{1} \geqslant 0$ in $\Sigma(\lambda)$ ). Thus

$$
\begin{equation*}
\Delta w+\Sigma b_{j} w_{j}+c w \leqslant 0 \tag{2.13}
\end{equation*}
$$

for suitable bounded functions $b_{j}, c$.
Note that (2.13) holds in $\Sigma(\lambda)$, for any $\lambda$ in ( $-a, 0$ ) provided we know $u_{1} \geqslant 0$ in $\Sigma(\lambda)$.

Since $w \geqslant 0$ in $\Sigma(\lambda)$ for $0<a+\lambda$ small, we infer from the maximum principle and the Hopf lemma that (2.10) and (2.11) hold. In fact we have proved

LEMMA 2.1. Assume that for some $\lambda$ in $(-a, 0)$ we have

$$
\begin{equation*}
u_{1}(x) \geqslant 0, u(x) \leqq u\left(x^{\lambda}\right) \text { in } \Sigma(\lambda) \tag{2.14}
\end{equation*}
$$

Then (2.10), (2.11) hold, i.e.

$$
\begin{equation*}
u(x)<u\left(x^{\lambda}\right) \text { in } \Sigma(\lambda) \text { and } u_{1}(\lambda, y)>0 \text { for } y \in \omega \tag{2.15}
\end{equation*}
$$

Step (I) is finished. Turn to Step (II). Inequalities (2.14) hold for every $\lambda$ in a maximal interval $(-a, \mu]$ in $(-a, 0]$. We will prove that $\mu=0$. Suppose the contrary, that $\mu<0$. By Lemma 2.1 we have (2.15) for $\lambda=\mu$.

Since $\mu$ is maximal, only two situations are possible. Case 1 . There is a sequence of points $x^{i}$ with $u_{1}\left(x^{i}\right)<0$ and $x_{1}^{i} \nLeftarrow \mu$. Case 2 . There are sequences $\lambda^{i} \downarrow \mu$, and $x^{i} \in \Omega$ with $x_{1}^{i}<\lambda^{i}$, such that

$$
\begin{equation*}
u\left(x^{i}\right)>u\left(\left(x^{i}\right)^{\lambda^{i}}\right) \tag{2.16}
\end{equation*}
$$

Consider first case 1. For a suitable subsequence we have $x^{i} \rightarrow x$ on $T_{\mu}$. Because of (2.15), necessarily $x \in \partial \Omega$. We may suppose that the exterior unit normal to $\partial \Omega$ at $x$ is $e_{2}=(0,1,0, \ldots, 0)$. From (2.6) we have $u_{1} \geqslant 0$ at $x$; by continuity, $u_{1}=0$ at $x$. Moving from $x^{i}$ along a segment in the direction $e_{2}$, in a short distance we hit $\partial S_{a}$, where $u_{1} \geqslant 0$. Consequently at some point on that segment we must have $u_{12} \geqslant 0$. So $u_{12}(x) \geqslant 0$. The function $w$ defined in (2.9), for $\lambda=\mu$, has in $\overline{\Sigma(\mu)}$ a minimum, zero, at the corner point $x$. By Lemma $S$ of [4] we have

$$
\left(\partial_{1}+\partial_{2}\right) w<0 \text { or }\left(\partial_{1}+\partial_{2}\right)^{2} w>0 \text { at } x
$$

i.e. $-2 u_{1}<0$ or $-4 u_{12}>0$ at $x$. But both are impossible. Thus Case 1 cannot occur.

Consider case 2 . We may suppose that in $\Sigma\left(\lambda^{i}\right), w=w_{\chi i}$ assumes its minimum, which is negative, at $x^{i}$. So $\nabla w=0$ and the Hessian matrix of spatial second derivatives of $w,\left\{w_{j k}\right\} \geqslant 0$ there. For a subsequence we have $x^{i} \rightarrow \bar{x} \in \overline{\Sigma(\mu)}$. In the limit $u(\bar{x}) \geqslant u\left(\bar{x}^{\mu}\right)$ so we must have equality. From now on $w$ refers to $w_{\mu}$. By continuity, at $x$,

$$
\begin{equation*}
w=0, \nabla w=0 \text { and } w_{j k} \geqslant 0 \tag{2.17}
\end{equation*}
$$

From (2.13) it follows that $\left\{w_{j k}\right\}=0$ there. By the theorem of the mean there is a point $z^{i}$ in the interval joining $x^{i}$ to $\left(x^{i}\right)^{\lambda^{i}}$ where $u_{1}\left(z^{i}\right)<0$. In view of (2.15) and the fact that Case 1 is impossible it follows that $\bar{x} \notin T_{\mu}$. Thus $\bar{x}=\left(\bar{x}_{1}, \bar{y}\right)$ with $-a \leqslant \bar{x}_{1}<\mu$. In $\overline{\Sigma(\mu)}, w$ achieves its minimum, zero, at $\bar{x}$.

If $-a=\bar{x}_{1}, \bar{y} \in \omega$ then by Lemma $H, w_{1}(\bar{x})>0$, contradiction. Se we have $\bar{y} \in \partial \omega$ and $\bar{x}_{1}>a$ or $\bar{x}_{1}=a$. We may suppose that the exterior normal to $\partial \omega$ at $\bar{y}$ is $(1,0, \ldots, 0)$. Then in the first case, $\bar{x}_{1}>a$, Lemma H implies $w_{2}(\bar{x})<0$, contradiction. In the second case, $\bar{x}_{1}=a$, Lemma $S$ implies

$$
\left(\partial_{1}-\partial_{2}\right)^{2} w(\bar{x})>0
$$

again a contradiction.
We have proved that $\mu=0$. To complete the proof of Theorem 2.1 we must prove the last assertion. Suppose $u_{1}(0, y)=0$ for some $y \in \omega$. For $w$ defined in (2.9) with $\lambda=0$ we have $w \geqslant 0$ in $\Sigma(0)$, and, as before, it satisfies an inequality of the form (2.13). It follows from the maximum principle that either $w \equiv 0$ or $w>0$. In the latter case, by Lemma $H$, we would have $-2 t_{1}=w_{1}<0$ at $(0, y)$, a contradiction. So $w \equiv 0$, i.e. $u$ is symmetric in $x_{1}$. Theorem 2.1 is proved.

REMARK 2.2. If we knew that $u_{1}>0$ near $x_{1}=-a$ then Step (I) in the proof would be trivial.

As in [4], Theorem 2.1 yields immediately the

Proof of Theorem 1.1: By Theorem 2.1 we have $u_{1}>0$ for $x_{1}<0$ and

$$
u(x) \leqslant u\left(x^{\lambda}\right) \text { for } \lambda=0
$$

But if we reflect the problem about the plane $x_{1}=0$, i.e. replace $x_{1}$ by $-x_{1}$, we obtain the same equation. Thus we may conclude that

$$
u(x) \geqslant u\left(x^{\lambda}\right) \text { for } \lambda=0
$$

Hence $u$ is symmetric in $x_{1}$.

Next we will prove a monotonicity result for solutions of (2.2) in $S_{a}$ under Neumann boundary conditions on $\partial \omega$.

THEOREM 2.2. Assume the conditions of Theorem 2.1 but with (2.6) replaced by

$$
\begin{equation*}
u_{\nu}\left(x_{1}, y\right)=0 \text { for } y \in \partial \omega \tag{2.6}
\end{equation*}
$$

where $\nu$ is the exterior unit normal to $S_{a}$ at $\left(x_{1}, y\right)$. Then for $-a<x_{1}<x_{1}^{\prime}$, $x_{1}+x_{1}^{\prime}<0$ and $y \in \bar{\omega}$ we have

$$
\begin{align*}
& u_{1}\left(x_{1}, y^{\prime}\right)>0  \tag{2.18}\\
& u\left(x_{1}, y\right)<u\left(x_{1}^{\prime}, y\right) \tag{2.19}
\end{align*}
$$

Furthermore if $u_{1}(0, y)=0$ for some $y \in \bar{\omega}$ then $u$ is a symmetric function of $x_{1}$.

The proof will rely on an extension of Proposition 1.1. in the appendix, Pro-
position A. 1.

Proof: (a) We proceed as in the proof of Theorem 2.1 - with slight variations. Using (2.4) we first show that (2.10) and (2.11) hold for $0<\lambda+a$ small. When using (2.4) we used Proposition 2.1. This was applied in the narrow region $-a<$ $<x_{1}<\lambda<\lambda_{0}, r \in \omega$; we had $w \geqslant 0$ on its boundary (except on top) and concluded that $w \geqslant 0$ in $\Sigma(\lambda)$. In the present situation we have $w \geqslant 0$ on the left $\left(x_{1}=\right.$ $=-a$ ) and right boundaries of $\Sigma(\lambda)$ and $\partial_{\nu} w=0$ on the remaining boundary for $\lambda<\lambda_{0}$. With the aid of Proposition A. 1 in the appendix we find $w \geqslant 0$ in $\Sigma(\lambda)$ for $\lambda$ close to $-a$. As before, using (2.13) we then find that $w>0$ in $\Sigma(\lambda)$.
(b) We have

LEMMA 2.2. L'nder the conditions of Lemma 2.1 we have (2.15) and

$$
\begin{equation*}
u_{1}(\lambda, y)>0 \text { for } y \in \bar{\omega} \tag{2.20}
\end{equation*}
$$

Proof: As before, (2.20) holds for $y \in \omega$. Suppose $u_{1}(\lambda, y)=0$ for some $y \in$ $\in \partial \omega$. We may suppose the exterior normal $\nu$ there is $e_{2}$. Applying Lemma $S$ to $w_{\lambda}$ in $\Sigma(\lambda)$ at $(\lambda, y)$ we infer that

$$
-4 u_{12}=\left(\partial_{1}+\partial_{2}\right)^{2} w>0 \text { there }
$$

But since $u_{2}=u_{\nu} \equiv 0$ on a segment containing ( $\lambda, y$ ) parallel to the $x_{1}$ axis we have $u_{12}(\lambda, y)=0$. Contradiction.
(c) With $\mu$ as in the proof of Theorem 2.1 we wish to show $\mu=0$. Suppose $\mu<0$. As before, we have the Cases 1 and 2. By Lemma $\tilde{2.2}$, Case 1 cannot occur so we have case $2: x^{i} \rightarrow \bar{x} \in \overline{\Sigma(\mu)}$; at $\bar{x}, w=0, \nabla w=0$ and $\left\{w_{j k}\right\}=0$ for $w=w_{\mu}$. So $\bar{x} \in \partial \Sigma(\mu)$, and in view of (2.20), $-a \leqslant \bar{x}_{1}<\mu$. If $-a<\bar{x}_{1}<\mu$ then $\bar{y} \in \partial \omega$ and by Lemma H we find $\partial_{\nu} w(\bar{x})<0$ - contradiction. If $-a=\bar{x}_{1}$, and $y \in \omega$ we see again by Lemma $H$ that $w_{1}(\bar{x})>0$ - contradiction. So $-a=$ $=\bar{x}_{1}, \bar{y} \in \partial \omega$. Using Lemma $S$ there we obtain again a contradiction. Thus $\mu=0$.

The last assertion of Theorem 2.1 holds as before. Theorem 2.2 is proved.

Theorems 2.1 and 2.2 have immediate applications to infinite cylinders. Let $S$ be the infinite cylinder $(-\infty, \infty) \times \omega$.

COROLLARY 2.1. Let $u$ be a $C^{2}$ solution of (2.2) in $\bar{S}$ with $f$ satisfying condition (2.3). Assume that for some numbers $c_{0}<c_{1}$, and $R>0$,

$$
\left\{\begin{array}{l}
\overline{\lim }_{x_{1} \rightarrow-\infty} u\left(x_{1}, y\right) \leqslant c_{0} \text { uniformly for } y \in \bar{\omega}  \tag{2.21}\\
u_{1}\left(x_{1}, y\right)>0 \text { for } x_{1}<-R \text { in } S \\
u\left(x_{1}, y\right) \geqslant c_{1} \text { for } x_{1} \geqslant-R \text { in } S .
\end{array}\right.
$$

(a) If everywhere on $\partial S$ we have $u_{1} \geqslant 0$, then $u_{1}>0$ in $S$.
(b) If everywhere on $\partial S$ we have $u_{\nu}=0$ then $u_{1}>0$ in $\bar{S}$.

Proof: Simply apply Theorems 2.1 and 2.2 and Remark 2.2, Lemma 2.2, in cylinders $\left\{-a<x_{1}<b\right\} \times \omega$ for arbitrarily large $a$ and $b$.

Extensions and applications of this kind of result are presented in our forthcoming paper [9]. In the next section we consider noncylindrical domains.

## 3. GENERAL SMOOTH DOMAINS

Here we extend the monotonicity results to more general domains. In addition, we present results on monotonicity in $x_{1}$ in the entire domain; see also the next section.

We will begin with an extension of Theorem 2.1 to general domains $\Omega$. For convenience we will assume $\Omega$ has no corners: $\Omega$ is a bounded domain with smooth boundary and convex in the $x_{1}$ direction. We suppose

$$
-a=\min \left\{x_{1} ; x \in \bar{\Omega}\right\} .
$$

Assume that

$$
\begin{equation*}
\left(x_{1}, y\right) \in \Omega, x_{1}<0 \Rightarrow\left(x_{1}^{\prime}, y\right) \in \Omega \text { for } x_{1}<x_{1}^{\prime}<-x_{1} . \tag{3.1}
\end{equation*}
$$

This implies that for $x \in \partial \Omega, x_{1}<0$, the unit exterior normal $\nu$ to $\partial \Omega$ at $x$ has $\nu_{1} \leqslant 0$.

In $\bar{\Omega}$ let $u$ be a $C^{3}$ solution of (2.2) where $f$ satisfies the conditions of Theorem 2.1. In particular

$$
\begin{align*}
& \text { for } x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0 \\
& f\left(x_{1}, y, u, p_{1}, p_{2}, \ldots, p_{n}\right) \leqslant f\left(x_{1}^{\prime}, y, u,-p_{1}, p_{2}, \ldots, p_{n}\right) \text { if } p_{1}>0 . \tag{3.2}
\end{align*}
$$

Corresponding to (2.5), (2.6) we make the following assumptions:

$$
\begin{align*}
& \text { If }\left(x_{1}, y\right) \in \partial \Omega \text { and }\left(x_{1}^{\prime}, y\right) \in \Omega, x_{1}<x_{1}^{\prime} \text { then } u\left(x_{1}, y\right)<u\left(x_{1}^{\prime}, y\right)  \tag{3.3}\\
& \left\{\begin{array}{l}
\text { If }\left(x_{1}, y\right),\left(x_{1}^{\prime}, y\right) \in \partial \Omega, x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0 \text { then } \\
u\left(x_{1}, y\right) \leqslant u\left(x_{1}^{\prime}, y\right) .
\end{array}\right.
\end{align*}
$$

From conditions (3.3), (3.4) it follows that

$$
\begin{equation*}
u_{1} \geqslant 0 \text { on } \partial \Omega \cap\left\{x_{1} \leqslant 0\right\} . \tag{3.5}
\end{equation*}
$$

As before we define for $\lambda>-a$,

$$
\begin{aligned}
& \Sigma(\lambda)=\left\{x \in \Omega ; x_{1}<\lambda\right\} \\
& \Sigma^{\prime}(\lambda)=\text { the reflection of } \Sigma(\lambda) \text { in the plane } T_{\lambda}=\left\{x_{1}=\lambda\right\}
\end{aligned}
$$

THEOREM 3.1. Assume the conditions above, i.e., (3.1)-(3.4). Then if $\left(x_{1}, y\right)$, $\left(x_{1}^{\prime}, y\right) \in \Omega, x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<0$ we have

$$
\begin{equation*}
u_{1}\left(x_{1}, y\right)>0, u\left(x_{1}, y\right)<u\left(x_{1}^{\prime}, y\right) \tag{3.6}
\end{equation*}
$$

Furthermore if $u_{1}=0$ for some points on $T_{0} \cap \Omega$ then $u$ is a symmetric function of $x_{1}$.

Proof: (i) $\ln \Sigma(\lambda)$ we consider, as before, the functions $v(x)=u\left(x^{\lambda}\right)$ and $w=w_{\lambda}$ given by (2.9), and we wish to show that for $-a<\lambda<0$, inequalities (2.10) and (2.11) hold. As before, using (3.2) one derives (2.12), and establishes (2.10), (2.11) for $\lambda$ close to $-a$. Thus if (2.10), (2.11) hold for all $\lambda$ in $\left(-a, \lambda_{0}\right)$, then (2.7) holds for $x_{1}<\lambda_{0}, x_{1}<x_{1}^{\prime}, x_{1}+x_{1}^{\prime}<2 \lambda_{0}$,
(ii) Next we have

LEMMA 3.1. Assume that for some $\lambda$ in $-a<\lambda<0$,

$$
\begin{equation*}
u_{1}(x) \geqslant 0, u(x) \leqslant u\left(x^{\lambda}\right) \text { in } \Sigma(\lambda) \tag{3.7}
\end{equation*}
$$

Then

$$
\left\{\begin{array}{l}
u_{1}(x)>0 \text { for } x \in T_{\lambda} \cap \Omega  \tag{3.8}\\
u(x)<u\left(x^{\lambda}\right) \text { in } \Sigma(\lambda)
\end{array}\right.
$$

Furthermore if $(\lambda, y) \in \partial \Omega$ and $\nu_{1}>-1 / 2$ there, then in a neighbourhood of $(\lambda, y)$ in $\Omega$ we have

$$
\begin{equation*}
u_{1}>0 \tag{3.9}
\end{equation*}
$$

Proof: The proof of (3.8) is similar to that of (2.15) and will be omitted. Turn to (3.9); at ( $\lambda, y$ ) we may suppose the unit normal $\nu=\left(\nu_{1}, \nu_{2}, 0, \ldots, 0\right)$. Since $\nu_{1}>-1 / 2$ we see that at $(\lambda, y)$ the two bounding surfaces of $\Sigma(\lambda), \partial \Omega$ and $T_{\lambda}$, meet at an angle $\theta>\pi / 3$ ( $\theta$ is the angle in $\Sigma(\lambda)$ ).

In $\Sigma(\lambda)$ the function $w=w_{\lambda}$ is positive and satisfies (2.13). At $(\lambda, y)$ it achieves its minimum. We may apply Lemma A. 1 of [4] and infer that at $(\lambda, y)$

$$
\left(\partial_{1}+b \partial_{2}\right) w<0 \text { or }\left(\partial_{1}+b \partial_{2}\right)^{2} w>0
$$

for $b$ large, i.e.

$$
u_{1}>0 \text { or } u_{12}<0 \text { at }(\lambda, y)
$$

If $u_{1}>0$ at $(\lambda, y)$ then we are through. Il $u_{12}<0$ at $(\lambda, y)$ then it is $<0$ in a neighbourhood in $\Omega$. But $u_{1} \geqslant 0$ on $\partial \Omega$. It follows that $u_{1}>0$ in a neighbourhood of $(\lambda, y)$ in $\Omega$.

The lemma is proved.
(iii) As before we have (3.7) for all $\lambda$ in a maximal interval $(-a, \mu]$ in $(-a, 0]$, and we wish to prove $\mu=0$. Suppose $\mu<0$. By Lemma 3.1, inequalities (3.8) hold.

We must treat the same two cases as in the proof of Theorem 2.1. Consider first Case 1: we have a sequence $x^{i}$ with $x_{1}^{i} \nLeftarrow \mu$ and $u_{1}\left(x^{i}\right)<0$. We will show that for $\epsilon>0$ sufficiently small, and for every $\lambda$ in some interval $\mu<\lambda<\lambda_{0}$, the function $w_{\lambda} \geqslant 0$ in the region

$$
\Omega_{\epsilon, \lambda}=\left\{x \in \Omega ; \mu-\epsilon<x_{1}<\lambda\right\} .
$$

This implies that $u_{1} \geqslant 0$ in $\Omega_{\epsilon, \lambda_{0}}$ - contradicting the assumption in case 1 .
To show $w_{\lambda}(x) \geqslant 0$ in the region

$$
\tilde{\Omega}_{\epsilon}=\left\{(x, \lambda) ; x \in \Omega, \mu-\epsilon<x_{1}<\lambda, \mu<\lambda<\lambda_{0}\right\}
$$

in ( $x, \lambda$ ) space, we will use the parabolic inequality (2.12), which follows, as before, using (3.2). We have to check that $w_{\lambda} \geqslant 0$ at every point on $\partial \tilde{\Omega}_{\epsilon}$ except those on $\lambda=\lambda_{0}$. As in $\S 2$ it will then follow with the aid of Proposition 1.1 that $w_{\lambda} \geqslant 0$ in $\widetilde{\Omega}_{\epsilon}$. We will show that

$$
\begin{equation*}
w_{\lambda} \geqslant 0 \text { on } \partial \Omega_{\epsilon, \lambda} \backslash T_{\lambda} \tag{3.10}
\end{equation*}
$$

In order to establish (3.10) we divide the points on $\partial \Omega_{\epsilon, \lambda} \backslash T_{\lambda}$ into different classes.

The set of points $K$ on $T_{\mu} \cap \partial \Omega$ at which $\nu_{1} \geqslant-1 / 3$ is compact. Thus for some $\delta>0$ we see from (3.9) that $u_{1} \geqslant 0$ at every point in $K_{\delta}$. Here

$$
K_{\rho}:=\{x \in \bar{\Omega} ; \operatorname{dist}(x, K)<\rho\}
$$

For $0<\epsilon, \lambda-\mu$ small, at any point

$$
\left(x_{1}, y\right) \in J:=\left\{x \in \partial \Omega \backslash K_{\delta / 2} ; \mu-\epsilon<x_{1}<\lambda\right\}
$$

we have $\nu_{1} \leqslant-1 / 4$ and $u\left(x_{1}^{\prime}, y\right)>u\left(x_{1}, y\right)$, if $\left(x_{1}^{\prime}, y\right) \in \Omega$ and $x_{1}^{\prime}>x_{1}-$ by (3.3). In particular for $\lambda-\mu>0$ but small,

$$
\begin{equation*}
w_{\lambda}(x)>0 \text { for } x \in J \tag{3.11}
\end{equation*}
$$

Consider the compact set

$$
L=\left(T_{\mu-\epsilon} \cap \bar{\Omega}\right) \backslash K_{\delta / 2}
$$

On $L$ we have $w_{\mu} \geqslant c_{1}$, a small positive constant. Hence for $0<\lambda-\mu$ small we also have

$$
\begin{equation*}
w_{\lambda}>0 \text { on } L \tag{3.11}
\end{equation*}
$$

We have established $w_{\lambda}>0$ at all points of $\partial \Omega_{\epsilon, \lambda} \backslash T_{\lambda}$ except those in $K_{\delta / 2}$. But for $0<\epsilon, \lambda-\mu$ small, we see that if $x \in K_{\delta / 2}$ and $\mu-\epsilon<x_{1}<\lambda$, then the interval joining $x$ to $x^{\lambda}$ lies in $K_{\delta}$ and so $u_{1} \geqslant 0$ on it. Hence $w_{\lambda}(x) \geqslant 0$. We have thus verified (3.10). Consequently $w_{\lambda} \geqslant 0$ in $\Omega_{\epsilon, \lambda}$ for $\mu<\lambda<\lambda_{0}$. This implies $u_{1} \geqslant 0$ in $\Omega_{\epsilon, \lambda_{0}}$. Contradiction.

Consider now case 2 . We have $\lambda^{i} \downarrow \mu$, and a sequence $x^{i} \rightarrow \bar{x} \in \overline{\Sigma(\mu)}$ with $u\left(x^{i}\right)>u\left(\left(x^{i}\right)^{\lambda i}\right)$, and $\nabla w_{\lambda i}=0$ at $x^{i}$. As before, since case 1 is impossible, $\bar{x} \in \partial \Omega$ and $\bar{x}_{1}<\mu$. Also $w_{\mu}=0, \nabla w_{\mu}=0$ at $\bar{x}$. By Lemma $H, \partial_{\nu} w_{\mu}(\bar{x})<0$, for $\nu$ the exterior unit normal at $\bar{x}$-contradiction.

Thus $\mu=0$ and the proof of Theorem 3.1 is complete except for the last assertion. That is proved just as in Theorem 2.1.

THEOREM 3.1'. Assume the conditions of Theorem 3.1 and assume that $\Omega$ is symmetric in $x_{1}$ and that the boundary values of $u$ are symmetric in $x_{1}$. Assume also $f$ is symmetric in $\left(x_{1}, p_{1}\right)$. Then $u$ is symmetric in $x_{1}$.

It is reasonable to ask if one can prove a monotonicity (in $x_{1}$ ) result in all of $\Omega$. Here is such a result. For convenience we suppose again that $\partial \Omega$ is smooth.

THEOREM 3.2. Let $\Omega$ be a bounded domain in $R^{n}$ with smooth boundary and which is convex in the $x_{1}$ direction. In $\Omega$ let $u$ be a $C^{3}$ solution of $(2.2)$ where $f$ satisfies (3.2) $f\left(x_{1}, y, u, p_{1}, \ldots, p_{n}\right) \leqslant f\left(x_{1}^{\prime}, y, u,-p_{1}, p_{2}, \ldots, p_{n}\right)$ if $p_{1} \geqslant 0$ for all $\left(x_{1}, y\right),\left(x_{1}^{\prime}, y\right)$ in $\Omega$ with $x_{1}<x_{1}^{\prime}$. Assume that on every interval in $\Omega$ parallel to the $x_{1}$-axis, with end points $x^{0}, x^{0}+t e_{1}$ on $\partial \Omega, t>0$, we have

$$
\begin{equation*}
u\left(x^{0}\right)<u\left(x^{0}+s e_{1}\right)<u\left(x^{0}+t e_{1}\right) \quad 0<s<t \tag{3.12}
\end{equation*}
$$

Assume also ( $A$ ): If $\partial \Omega$ contains a segment parallel to the $x_{1}$ axis on which $u$ is constant then $\nu$, the unit exterior normal to $\partial \Omega$, is also constant on the segment. (This condition is automatic if $n=2$ ). Assume also that at any boundary point of $\Omega$, where $\nu_{1}=0$, we have

$$
\begin{equation*}
u_{1} \geqslant 0 \tag{3.12}
\end{equation*}
$$

Then

$$
\begin{equation*}
u_{1}>0 \text { in } \Omega \tag{3.13}
\end{equation*}
$$

From (3.12) and (3.12), it follows that $u_{1} \geqslant 0$ everywhere on $\partial \Omega$ - so one might expect (3.13) to follow with the aid of the maximum principle. But the principle is not immediately applicable.

Proof: We suppose that $\min x_{1}$ and $\max x_{1}$ in $\bar{\Omega}$ are $-a$ and 0 , and we define $\Sigma(\lambda)$ and its reflection $\Sigma^{\prime}(\lambda)$ as above. For $\lambda$ greater than but close to $-a$ we have $\Sigma^{\prime}(\lambda) \subset \Omega$, but this will no longer be the case as we increase $\lambda$. In place of $\Sigma(\lambda)$ we will work with

$$
\Omega(\lambda)=\left\{x \in \Sigma(\lambda) ; x^{\lambda} \in \Omega\right\}
$$

In $\Omega(\lambda)$ we consider the function $w_{\lambda}(x)$ of (2.9). From our conditions it follows that $w=w_{\lambda} \geqq 0$ on $\partial \Omega(\lambda)$. To prove'the theorem we will prove the analogues of (2.10), (2.11): for $-a<\lambda<0$.

$$
\begin{equation*}
w_{\lambda}(x)>0 \text { in } \Omega(\lambda) \tag{3.14}
\end{equation*}
$$

$$
\begin{equation*}
w_{1}<0 \text { on } T_{\lambda} \cap \Omega \tag{3.15}
\end{equation*}
$$

The arguments will be similar to those used above.
(i) Just as before we find that (3.14) and (3.15) hold for $\lambda$ close to $-a$.
(ii) The analogue of Lemma 3.1 (and proved in the same way) is

LEMMA 3.1'. Assume that for some $\lambda$ in $-a<\lambda<0$,

$$
\begin{equation*}
u_{1}(x) \geqslant 0, u(x) \leqslant u\left(x^{\lambda}\right) \text { in } \Omega(\lambda) \tag{3.16}
\end{equation*}
$$

Then

$$
\begin{align*}
& u_{1}(x)>0 \text { for } x \in T_{\lambda} \cap \Omega \\
& u(x)<u\left(x^{\lambda}\right) \text { in } \Omega(\lambda) \tag{3.17}
\end{align*}
$$

Furthermore if $(\lambda, y) \in \partial \Omega$ and $\left|\nu_{1}\right|<1 / 2$ there, then in a neighbourhood of $(\lambda, y)$ in $\Omega$ we have

$$
\begin{equation*}
u_{1}>0 \tag{3.18}
\end{equation*}
$$

(iii) We have (3.16) for all $\lambda$ in a maximal interval ( $-a, \mu$ ] in ( $-a, 0$ ], and we wish to prove $\mu=0$. As before, suppose $\mu<0$. By Lemma $3.1^{\prime}$, inequalities (3.17) hold for $\lambda=\mu$. As usual we have two cases to exclude.

Case 1: There is a sequence $x^{i}$ with $x_{1}^{i} \ngtr \mu$ and $u_{1}\left(x^{i}\right)<0$. We will show that for $\epsilon>0$, sufficiently small, and for every $\lambda$ in some interval $\mu<\lambda<\lambda_{0}$ the function $w_{\lambda} \geqslant 0$ in the region (which may have infinitely many components):

$$
\Omega(\epsilon, \lambda)=\left\{x \in \Omega(\lambda) ; \mu-\epsilon<x_{1}<\lambda\right\},
$$

This implies that $u_{1} \geqslant 0$ in $U_{\mu \leqslant \lambda<\lambda_{0}} \Omega(\epsilon, \lambda)$, which contains a neighbourhood of $T_{\mu} \cap \Omega$, so that case 1 cannot hold.

As in the proof of Theorem 3.1 we wish to verify that

$$
\begin{equation*}
w_{\lambda} \geqslant 0 \text { on } \partial \Omega(\epsilon, \lambda) . \tag{3.19}
\end{equation*}
$$

On $T_{\lambda}$ we have $w_{\lambda}=0$. As before we divide the points on $\partial \Omega(\epsilon, \lambda) \backslash T_{\lambda}$ into different classes.

The set $K$ of points on $T_{\mu} \cap \partial \Omega$ at which $|\nu| \leqslant 1 / 3$, is compact. So for some $\delta>0$ we see from (3.9) that $u_{1} \geqslant 0$ at every point in $K_{\delta}$. For $\epsilon, \lambda-\mu$ small, on

$$
J:=\left\{x \in \partial \Omega \backslash K_{\delta / 2} ; \mu-\epsilon<x_{1}<\lambda \text { and } \nu_{1}<0\right\}
$$

we have $\nu_{1} \leqslant-1 / 4$. If $\left(x_{1}, y\right) \in J,\left(x_{1}^{\prime}, y\right) \in \Omega, x_{1}^{\prime}>x_{1}$ then $u\left(x_{1}^{\prime}, y\right)>u\left(x_{1}, y\right)$.
Thus $w_{\lambda}(x)>0$ for $x \in J$. Let $M$ be the set of points $x \in\left(\partial \Omega(\epsilon, \lambda) \backslash K_{\delta / 2}\right)$ with $\mu-\epsilon<x_{1}<\lambda$ such that $x^{\lambda} \in \partial \Omega$. Then $\nu_{1}\left(x^{\lambda}\right) \geqslant 1 / 4$ and we find from (3.12) that

$$
w_{\lambda}(x)>0 \text { for } x \in M .
$$

On the compact set

$$
L=\left(T_{\mu-\epsilon} \cap \partial \Omega(\epsilon, \lambda)\right) \backslash K_{\delta / 2}
$$

we have $w_{\mu}>c_{1}>0$. So for $0<\lambda-\mu$ small we also have $w_{\lambda}>0$ on $L$. Thus $w_{\lambda}>0$ at all points of $\partial \Omega(\epsilon, \lambda) \backslash T_{\lambda}$ except possibly those in $K_{\delta / 2}$. But as before we see that $w_{\lambda}(x) \geqslant 0$ there. So (3.19) is verified and case 1 is impossible.

Case 2. There is a sequence $\lambda^{i} \downarrow \mu$ and a sequence of points $x^{i} \in \Omega\left(\lambda^{i}\right)$, with $u\left(x^{i}\right)>u\left(\left(x^{i}\right)^{\lambda^{l}}\right)$. So $u(\bar{x}) \geqslant u\left(\bar{x}^{\mu}\right)$. As before, since case 1 is impossible, it follows that $\bar{x}_{1}<\mu$ and $\bar{x} \in \overline{\Omega(\mu)}$. Also $w_{\mu}(\bar{x})=0$. It follows from (3.12) and (3.12)' that the straight segment $\sigma$ joining $\bar{x}$ and $\bar{x}^{\mu}$ belongs entirely to $\partial \Omega$, and $u=$ constant on $\sigma$. Now if $n>2$, at $\bar{x}$, the boundary of $\Omega(\mu)$ may have a sharp corner. So we cannot apply Lemma $H$. For that reason we assumed condition (A). Because of that condition, we may apply Lemma H in $\Omega(\mu)$ at $\bar{x}$, and conclude that $\partial_{\nu} w<0$. We may then proceed as in the proof of Theorem 3.1 and we see that Case 2 is impossible. Thus $\mu=0$ and the proof of Theorem 3.2 is complete.

REMARK 3.1. In the theorem, if $\Omega$ is symmetric in $x_{1}$ about $x_{1}=\bar{a} / 2$ then Hypothesis (A) is automatically satisfied.

Can we drop Hypothesis (A) in general?

A monotonicity result in a full cylinder $\Omega=S_{a}$ may be proved in a similar way. The details are simpler and we merely state the result.

THEOREM 3.3. Let $u$ be a $C^{2}$ solution in $\bar{\Omega}$ of (2.2) where $f$ satisfies (3.2) for $x_{1}<x_{1}^{\prime}$. Assume

$$
\left\{\begin{array}{l}
u(-a, y) \leqslant u\left(x_{1}, y\right) \leqslant u(a, y) \text { for }-a<x_{1}<a, y \in \omega,  \tag{3.20}\\
\text { and } \forall x_{1} \text { in }-a<x_{1}<a, \exists y, y^{\prime} \in \omega \text { such that. } \\
u(-a, y)<u\left(x_{1}, y\right), u\left(x_{1}, y^{\prime}\right)<u\left(a, y^{\prime}\right)
\end{array}\right.
$$

Assume also

$$
\begin{equation*}
u_{1}\left(x_{1}, y\right) \geqslant 0 \text { for } y \in \partial \omega \tag{3.21}
\end{equation*}
$$

Then

$$
\begin{equation*}
u_{1}>0 \text { in } \Omega \tag{3.22}
\end{equation*}
$$

Berestycki and Pacella [2] adapted the method of moving planes to derive symmetry results in sector-like domains. Here too we may consider other geometries and derive other forms of Theorem 3.2. One may use the method of moving planes but not require them to be parallel while moving. As an example we present such a result in an angular sector.

In $R^{n}$ let $\left(\rho, \theta_{1}, \ldots, \theta_{n-1}\right)$ be polar coordinates, $\rho \geqslant 0, \theta_{i} \in[0, \pi]$ for $1 \leqslant$ $\leqslant i \leqslant n-2, \theta_{n-1} \in[0,2 \pi)$. In $R^{n}$ let $\Omega$ be a bounded domain whose closure does not touch the $x_{n}$-axis and which, in polar coordinates, is a product domain

$$
\omega \times\left(0<\theta_{n-1}<\alpha\right)
$$

Here $\omega$ is a bounded domain in $R^{n-1}$ with smooth boundary, and $\alpha<2 \pi$.
THEOREM 3.4. Let $u \in C^{2}(\bar{\Omega})$ be a solution of

$$
\Delta u+f\left(\rho, \theta_{1}, \ldots, \theta_{n-1}, u, u_{\rho}, u_{\theta_{1}}, \ldots, u_{\theta_{n-1}}\right)=0
$$

with $f$ continuous, and Lipschitz continuous in $(u, \nabla u)$, Assume (for convenience) the analogue of (3.2) (here $p=\left(p_{1}, \ldots, p_{n-1}\right)$ ).

$$
\begin{align*}
& f\left(\rho, \theta_{1}, \theta_{2}, \ldots, \theta_{n-1}, u, u_{\rho}, p_{1}, \ldots, p_{n-1}\right) \leqslant \\
& \leqslant f\left(\rho, \theta_{1} \ldots, \theta_{n-1}^{\prime}, \theta_{n}, u, u_{\rho}, p_{1}, p_{2}, \ldots,-p_{n-1}\right)  \tag{3.23}\\
& \text { if } \theta_{n-1}<\theta_{n-1}^{\prime}, p_{n-1} \geqslant 0
\end{align*}
$$

Assume that for every fixed $\sigma=\left(\rho, \theta_{1}, \ldots, \theta_{n-2}\right)$ (corresponding to a point in $\omega$ ),

$$
\left\{\begin{array}{l}
u(\sigma, 0) \leqslant u(\sigma, s) \leqslant u(\sigma, \alpha) \text { for } 0<s<\alpha  \tag{3.24}\\
\text { and } \forall s \text { in } 0<s<\alpha, \exists \sigma, \sigma^{\prime} \text { such that } u(\sigma, 0)<u(\sigma, s) \\
u\left(\sigma^{\prime}, s\right)<u\left(\sigma^{\prime}, \alpha\right)
\end{array}\right.
$$

Assume also $u_{\theta_{n-1}} \geqslant 0$ on $\partial \Omega \cap\left\{0<\theta_{n-1}<\alpha\right\}$. Then $u_{\theta_{n-1}}>0$ in $\Omega$.
The proof is, again, left to the reader.
L. Caffarelli has pointed out to us a much simpler argument to prove monotonicity in a full cylinder - and even uniqueness for solutions satisfying (roughly) (3.24). Here is such a result, and his argument .

THEOREM 3.5. In the cylinder $\Omega=S_{a}$ let u be a $C^{2}(\Omega) \cap C^{1}(\bar{\Omega})$ solution of (2.2) with $f$ continuous and Lipschitz continuous in $(u, \nabla u)$. Assume that $f$ is nondecreasing in $x_{1}$. Assume

$$
\begin{equation*}
u(-a, y)<u\left(x_{1}, y\right)<u(a, y) \text { for }-a<x_{1}<a, y \in \omega \tag{3.20}
\end{equation*}
$$

and either

$$
\begin{equation*}
u \text { is strictly increasing in } x_{1}, \text { on } \partial \Omega \tag{3.21}
\end{equation*}
$$

or the Neumann condition

$$
\begin{equation*}
u_{v}=0 \text { for }-a<x_{1}<a, y \in \partial \omega \tag{3.21}
\end{equation*}
$$

Then (a) $u_{1} \geqslant 0$ in $\Omega$. (b) If $f$ is also Lipschitz in $x_{1}$ then $u_{1}>0$ in $\Omega$. (c) There is at most one solution with given Dirichlet data, or with Neumann conditions $(3.21)^{\prime \prime}$, satisfying all these conditions, i.e. u is unique.

Proof: (a) In the region $\Sigma(\lambda)$, for $-a<\lambda<a$, consider the function $v(x)=$ $=u\left(x_{1}+a-\lambda, y\right)$, and set $w=v(x)-u(x)$. $w$ satisfies

$$
\begin{aligned}
& 0=\Delta w+f\left(x_{1}+a-\lambda, y, v, \nabla v\right)-f(x, u, \nabla u)=0 \\
& \geqslant \Delta w+f(x, v, \nabla v)-f(x, u, \nabla u) \\
& =\Delta w+\Sigma b_{j} w_{j}+c w
\end{aligned}
$$

for suitable coefficients $b_{j}, c$, having $L^{\infty}$ norms bounded by a constant independent of $\lambda$. Also $w$ 짖 0 on $\partial \Sigma(\lambda)$, or in case of (3.21)", on part of $\partial \Sigma(\lambda)$, with $w_{v}=0$ on the remaining part. We use the «sliding domain method»: For $0<a+$ $+\lambda$ small we infer that $w>0$ in $\Sigma(\lambda)$, i.e. $v$ lies above $u$. Now increase $\lambda$. On $\partial \Sigma(\lambda)$ we always have $w=w_{\lambda} \geq 0$ for $\lambda<a$. So by the maximum principle we must continue to have $w>0$ for every $\lambda<a$. This implies that $u$ is increasing in $x_{1}$.
(b) Differentiating the equation (2.2) with respect to $x_{1}$ and applying the maximum principle again we find $u_{1}>0$.
(c) If $\underline{u}$ is another solution satisfying the same conditions, the same argument shows that $v(x)=u\left(x_{1}+a-\lambda, y\right)$ is greater than $\underline{u}$ in $\Sigma(\lambda)$ for $0<a+\lambda$ small. Increasing $\lambda$ as before we find $\underline{u}(x) \geqslant u(x)$. Interchanging the roles of $u$ and $\underline{u}$ we find the opposite inequality.

COROLLARY 3.1. Let $u, f$ satisfy the conditions of Theorem 3.5. Assume in addition that the boundary values of $\phi$ are odd in $x_{1}$. Assume also that $f$ is odd in $\left(x_{1}, u, p_{2}, \ldots, p_{\mathrm{n}}\right)$. Then $u$ is odd in $x_{1}$.

This follows, from the fact that $\underline{u}=-u\left(-x_{1}, y\right)$ is a solution satisfying the same conditions, and so equals $u$.

REMARK 3.2. The conditions in Theorems 3.3 and 3.5 are slightly different. Neither implies the other. Here is an example of an equation where the conditions of Theorem 3.3 hold but not those of Theorem 3.5 . On the interval $-1 \leqslant x \leqslant 1$ the function $u=x$ satisfies the equation $\ddot{u}+\min (0, x) \dot{u}+\max (0, x) \dot{u}^{2}-u=0$.

REMARK 3.3. In Theorems 3.1, $3 . \overline{1}$ and 3.2 we assumed $u \in C^{3}(\bar{\Omega})$. Li, Cong Ming pointed out to us that with slight modifications the proofs work if $u \in C^{2, \alpha}$ ( $\bar{\Omega}$ ) for some $\alpha>0$. This is because in Lemma A.1 of [4], the last assertion holds if $u \in C^{r, \alpha}$ in $\bar{\Omega}$ near 0 for $r+\alpha>\pi / \theta_{0}$. See Lemma A. 1 in this paper.

The argument used in the proof of Theorem 3.5 yields the following result.
THEOREM 3.6. Let $f$ be as in Theorem 3.5 and let $u, \underline{u} \in C^{3}(\bar{\Omega})$ be solutions of (2.2) satisfying

$$
\begin{equation*}
\underline{u}\left(x_{1}, y\right)<u(a, y), \underline{u}(-a, y)<u\left(x_{1}, y\right) \tag{3.25}
\end{equation*}
$$

for $-a<x_{1}<a, y \in \omega$. Assume also either

$$
\begin{equation*}
\underline{u}\left(x_{1}, y\right)<u\left(x_{1}^{\prime}, y\right) \text { for }-a \leqslant x_{1}<x_{1}^{\prime} \leqslant a, y \in \partial \omega, \tag{3.26}
\end{equation*}
$$

or

$$
\begin{equation*}
\underline{u}_{v}=u_{\nu}=0 \text { for }-a<x_{1}<a, y \in \partial \omega . \tag{3.27}
\end{equation*}
$$

Then

$$
\underline{u} \leqslant u \text { in } \Omega .
$$

## 4. MONOTONICITY IN THE WHOLE DOMAIN AND ANTISYMMETRY

Theorems 3.3 and 3.5 concerned monotonicity in cylinders. Corollary 3.1 applied the latter to prove antisymmetry. In this section we will present a stronger version of Theorem 3.5 as well as Cor. 3.1.

Consider again the finite cylinder $\Omega=S_{a}=(-a, a) \times \omega$, with $a>0$; as before $\omega \subset R^{n-1}$ is a bounded domain with smooth boundary. Let $u \in C^{2}(\bar{\Omega})$ be a solution of

$$
\begin{align*}
& \Delta u+f(x, u, \nabla u)=0 \text { in } \Omega  \tag{4.1}\\
& u=\phi \text { on } \partial \Omega
\end{align*}
$$

Here $\phi$ is a continuous function on $\partial \Omega$ satisfying

$$
\begin{equation*}
\phi\left(x_{1}, y\right) \leqslant \phi\left(x_{1}^{\prime}, y\right) \text { for } x_{1} \leqslant x_{1}^{\prime} \tag{4.3}
\end{equation*}
$$

The function $f(x, y, p)$ is continuous in all variables, locally Lipschitz in ( $u, p$ ) and satisfies

$$
\begin{equation*}
f(x, u, p) \text { is nondecreasing in } x_{1} \text { for } p_{1} \geqslant 0 \tag{4.4}
\end{equation*}
$$

THEOREM 4.1. Let $u$ be as above and assume conditions (4.3), (4.4). Assume also that $u$ satisfies

$$
\left\{\begin{array}{l}
\phi(-a, y) \leqslant u\left(x_{1}, y\right) \leqslant \phi(a, y) \text { for }-a<x_{1}<a, y \in \omega  \tag{4.5}\\
\text { and } \forall x_{1} \text { in }(-a, a), \exists y \in \omega \text { such that } \\
\phi(-a, y)<u\left(x_{1}, y\right)
\end{array}\right.
$$

Then $u$ is strictly increasing in $\mathrm{x}_{1}$ in $\Omega$. Furthermore it is unique, i.e., if $\underline{u}$ is another solution of (4.1), (4.2) satisfying (4.5) then $\underline{u}=u$.

COROLLARY 4.1. Under the additional assumptions that $u_{1} \in C^{2}(\Omega)$ and that $f$ is locally Lipschitz in $\left(x_{1}, u, p\right)$ we have the stronger conclusion: $u_{1}>0$ in $\Omega$.

COROLLARY 4.2 (ANTISYMMETRY). Assume the conditions of Theorem 4.1 and assume in addition that $\phi$ is odd in $x_{1}$, on $\partial \Omega$, and that $f(x, u, p)$ is odd in $\left(x_{1}, u, p_{2}, \ldots, p_{n}\right)$. Then $u$ is odd, i.e. antisymmetric in $x_{1}$ :

$$
u\left(-x_{1}, y\right)=-u\left(x_{1}, y\right) \forall\left(x_{1}, y\right) \text { in } \Omega
$$

COROLLARY 4.3. Let $f$ be in Theorem 4.1 and satisfy, in addition, for constants $A<B$,

$$
f(x, u, 0) \geqslant 0 \quad \text { if } u \leqslant A, f(x, u, 0) \leqslant 0 \text { if } u \geqslant B
$$

Suppose $u$ and $\underline{u} \in C^{2}(\bar{\Omega})$ are solutions of (4.1), (4.2) where $\phi$ satisfies (4.3) and

$$
\phi \equiv A \quad \text { on } x_{1}=-a, \quad \phi \equiv B \text { on } x_{1}=a
$$

Then $u \equiv \underline{u}$, i.e. the solution is unique, and $u$ is strictly increasing in $x_{1}$ in $\Omega$.

Proof: This follows from the Theorem if we can show that

$$
A<u, \underline{u}<B \text { in } \Omega
$$

We will just show $u<B$; the other inequalities are proved in the same way. If $M=\max u>B$ then in a neighbourhood of a point where $u=M$ we have

$$
\Delta u+f(x, u, p)-f(x, u, 0) \geqslant 0
$$

Using the Lipschitz continuity of $f$ in $p$ we find in that neighbourhood,

$$
\Delta u+\Sigma b_{i} u_{i} \geqslant 0
$$

with $b_{i} \in L^{\infty}$. By the maximum principle, the set of points where $u=M$ is open. By continuity it is closed - and so all of $\bar{\Omega}$. Impossible. Next, if $u=B$ at some point in $\Omega$ then in $\Omega$ we have

$$
\Delta(u-B)+f(x, u, p)-f(x, B, 0) \geqslant 0
$$

Using Lipschitz continuity in ( $u, p$ ) we find

$$
\Delta(u-B)+b_{j}(u-B)_{j}+c(u-B) \geqslant 0 \text { in } \Omega
$$

Here the $b_{j}$ and $c$ are bounded measurable functions. But $u-B \leqslant 0$ in $\Omega$, with equality holding at some point in $\Omega$. The maximum principle implies $u \equiv B$. Impossible.

Even for $n=1$ the result seems new - even for the simple equation

$$
\ddot{u}+b(x) \dot{u}+f(u)=0
$$

with $b$ nondecreasing in $x$, and $g(u) \geqslant 0$ for $u \leqslant A, g(u) \leqslant 0$ for $u \geqslant B$. Li, CongMing has constructed an example with $b$ decreasing for which there is nonuniqueness.

REMARK 4.1. The condition (4.5) cannot be dropped. For example, in $R^{2}$, with $\omega=(0, \pi), \Omega=(-\pi / 2, \pi / 2) \times(0, \pi)$, the positive eigenfunction $\tilde{u}=\sin y \cos x_{1}$ of

$$
\begin{aligned}
& \Delta \tilde{u}+2 \tilde{u}=0 \text { in } \Omega \\
& \tilde{u}=0 \text { on } \partial \Omega
\end{aligned}
$$

is not antisymmetric in $x_{1}$. Using this example one may in fact obtain a more interesting one: a solution $u$ of the same equation which is not antisymmetric in $x_{1}$, and with boundary values $\phi$ which are strictly increasing (and odd) in $x_{1}$. Namely, let $v$ be the solution of

$$
\begin{aligned}
& \Delta v+2 v=0 \text { in } \Omega \\
& v=\widetilde{\phi} \text { on } \partial \Omega
\end{aligned}
$$

where $\bar{\phi}$ is odd and strictly increasing in $x_{1}-$ it is easily verified that this is always solvable. Then take $u=\tilde{u}+\epsilon u$.

The proof of Theorem 4.1 which we will present in an extension of that of Caffarelli's of Theorem 3.5. We have another proof of Cor. 4.2 which is somewhat surprising in that it makes use of the method of moving planes and reflection. However in the region $\Sigma(\lambda)$, in place of the function $w$ defined in (2.9), one works with the function $\widehat{w}=u(x)+u\left(x^{\lambda}\right)$. It satisfies Neumann boundary data on $x_{1}=\lambda$. One proves that $\tilde{w} \leqslant 0$ in $\Sigma(\lambda)$.

Suppose $\underline{u}$ and $u$ are solution satisfying all the conditions of the Theorem. We will show that in our usual $\Sigma(\lambda)$, for $-a<\lambda<a$, the (new) function

$$
\begin{equation*}
w(x):=u\left(x+(a-\lambda) e_{1}\right)-\underline{u}(x)>0 \tag{4.6}
\end{equation*}
$$

If we take $\underline{u}=u$ we infer that $u$ is strictly increasing in $x_{1}$. If we set $\lambda=a$ we see that $u \geqslant \underline{u}$ in $\Omega$. Interchanging the roles of $u$ and $\underline{u}$ it follows that $u=\underline{u}$. Thus (4.6) yields Theorem 4.1.

To prove (4.6) we will derive a parabolic inequality for $w$ of the form (2.12). But then we have need of parabolic analogues of Lemmas $H$ and S. So we present forms of these which will suffice for our purposes. We recall from section 1 that $V$ is a bounded domain in $R^{n+1}$ lying in $t<T$.

Hypotheses:

1) Here $w$ is a solution in $\bar{V}$ (see section 1) of (1.9):

$$
\begin{equation*}
\left(L-\beta \partial_{t}\right) w=a_{i j}(x, t) w_{x_{i} x_{j}}+b_{i}(x, t) w_{x_{i}}+c(x, t) w-\beta(x, t) w_{t} \geqslant 0 \tag{4.7}
\end{equation*}
$$

where the $a_{i j}$ are continuous and satisfy (1.8), and the other coefficients are in $L^{\infty}$, also $\beta \geqslant 0, w$ is supposed to have continuous second derivatives in the space variables $x$, and continuous time derivative $\partial_{t} w$ in $\bar{V}$.
2) $V_{\tau}=V \cap\{t=\tau\} \forall \tau<T$ is connected, and $V_{T}$ is connected. Here $V_{T}$ consists of points $(x, T)$ such that the lower open half of some ball with centre $(x, T)$ is in $V$. Set

$$
V \cup V_{T}=\tilde{V}
$$

In the following, $P$ denotes a paraboloid

$$
P=\left\{(x, t) ; t-T+\delta>\left|x-x^{0}\right|^{2}\right\}, \quad \delta>0
$$

for which the parabolic cap

$$
\begin{equation*}
Q=P \cap\{t \leqslant T\} \text { lies in } \tilde{V} \tag{4.8}
\end{equation*}
$$

We also consider parabolic caps with $T$ replaced by some other value.
Here is a parabolic analogue of Lemma H (for $\beta>0$ see [7], chapt. 3, section 3).
LEMMA 4.1. ( $\tilde{H})$ Let $V, w$ and $Q$ be as above. Suppose $w<0$ in $Q$ and equals zero at a point $(\bar{x}, T) \in \partial Q \cap \partial P$. Then $w_{\nu}>0$ there, where $\nu$ denotes any spatial outer direction to the sphere $\left|x-x^{0}\right|^{2}=\delta$ in the plane $t=T$.

Using this one easily establishes

LEMMA 4.2. Assume hypotheses 1), 2), and that $w \leqslant 0$ in $V$. Suppose $w<0$ at some point $\left(x^{0}, t^{0}\right) \in \tilde{V}$. Then $w<0$ on all of $V_{t}$.

Next an analogue of Lemma $S$.

LEMMA 4.3 ( $\widetilde{\mathbf{S}})$. Assume hypotheses 1 ), 2) with $w \leqslant 0$ in $\widetilde{V}$ and $w=0$ at a point $(\bar{x}, T)$ on $\partial V_{T}$. Suppose that near $(\bar{x}, T), \partial V \backslash V_{T}$ consists of two transversally intersecting $C^{2}$ hypersurfaces $\{\rho=0\}$ and $\{\sigma=0\}$, with $\rho, \sigma<0$ in $V$ and $\nabla_{x} \rho$, $\nabla_{x}$ olinearly independent at $(\bar{x}, T)$. Assume that at $(\bar{x}, T)$;

$$
\begin{equation*}
a_{i j} \rho_{x_{i}} \sigma_{x_{j}}=0 \tag{4.9}
\end{equation*}
$$

and assume that there exists a $C^{2}$ curve $\mathscr{C}$ of the form $(\xi(t), t)$ for $T-\epsilon<t \leqslant T$, lying in $\{\rho=\sigma=0\}$ such that on $\{\rho=\sigma=0\}$ we have

$$
\begin{equation*}
a_{i j} \rho_{x_{i}} \sigma_{x_{j}} \geqslant-C(\text { distance to } \mathscr{C})^{2} \tag{4.10}
\end{equation*}
$$

Conclusion: For any spatial outer direction $\nu$ at $(\bar{x}, T)\left(\right.$ i.e. $\left.\nabla_{x} \rho \cdot \nu, \nabla_{x} \sigma \cdot \nu>0\right)$ either

$$
\begin{equation*}
\partial_{\nu} w>0 \quad \text { or } \partial_{\nu}^{2} w<0 \text { at }(\bar{x}, T) \tag{4.11}
\end{equation*}
$$

We also have

LEMMA 4.4. Assume the conditions of Lemma 4.3 except that, in place of (4.9) we have

$$
\begin{equation*}
a_{i j} \rho_{x_{i}} \sigma_{x_{j}}>0 \text { at }(\bar{x}, T) \tag{4.9}
\end{equation*}
$$

Then

$$
\begin{equation*}
\partial_{v} w(\bar{x}, T)>0 \tag{4.12}
\end{equation*}
$$

For the convenience of the reader, proofs of these lemmas will be sketched in the Appendix. See also Lemma A. 1 there, a parabolic form of Lemma A. 1 of [4].

Proof of Theorem 4.1: As indicated earlier, the theorem follows once (4.6) is established, i.e.,

$$
\begin{equation*}
w(x, \lambda)=w(x)=v(x)-\underline{u}(x)>0 \text { in } \Sigma(\lambda), \text { for }-a<\lambda<a . \tag{4.6}
\end{equation*}
$$

Here

$$
\begin{equation*}
v(x)=u\left(x+(a-\lambda) e_{1}\right) \quad \text { in } \quad \Sigma(\lambda) \tag{4.13}
\end{equation*}
$$

We will prove (4.6)' by deriving a parabolic differential inequality for $w$ and then using the maximum principle. We have

$$
\begin{aligned}
& 0=\Delta w+f\left(x_{1}+a-\lambda, y, v, \nabla v\right)-f(x, \underline{u}, \nabla \underline{u}) \\
& =\Delta w+f(x, v, \nabla v)-f(x, \underline{u}, \nabla \underline{u}) \\
& +f\left(x_{1}+a-\lambda, y, v, \nabla v\right)-f(x, v, \nabla v)
\end{aligned}
$$

By condition (4.4) we see that

$$
I=f\left(x_{1}+a-\lambda, y, v, \nabla v\right)-f(x, v, \nabla v)
$$

satisfies $I \geqslant 0$, if $v_{1} \geqslant 0$ while, by Lipschitz continuity, $I \geqslant C v_{1}$ for some constant $C>0$, if $v_{1}<0$. Thus $I \geqslant \beta v_{1}$ where $\beta \geqslant 0$ is in $L^{\infty}$. Hence

$$
0 \geqslant \Delta w+f(x, v, \nabla v)-f(x, \underline{u}, \nabla \underline{u})+\beta v_{1}
$$

or

$$
\begin{equation*}
0 \geqslant \Delta w+b_{j} w_{j}+c w \cdots \beta \partial_{\lambda} w, \tag{4.14}
\end{equation*}
$$

by the integral theorem of the mean and the identity $v_{1}=-\partial_{\lambda} w$. This is our parabolic inequality. It holds in the region U in $(x, \lambda)$ space:

$$
U=\left\{(x, \lambda) ; x \in \Omega,-a<x_{1}<\lambda,-a<\lambda<a\right\}
$$

On the «spatial» boundary of $U$, i.e. the part of the boundary lying in $\lambda<a$ we have by (4.3) and (4.5), $w \geqslant 0$. In fact if we set

$$
\begin{equation*}
U_{t}=U \cap\{\lambda=t\}, \quad-a<t<a \tag{4.15}
\end{equation*}
$$

then

$$
\begin{equation*}
w \geqq 0 \text { on } \partial U_{t} . \tag{4.16}
\end{equation*}
$$

Note that $U_{t}=\Sigma(t)$ is connected.
For $0<a+\lambda$ small we may apply Proposition 1.1 and .infer that

$$
w(x, \lambda) \geqslant 0
$$

Since $U_{\lambda}=\Sigma(\lambda)$ is connected it follows from (4.16) and Lemma 4.2 that $w(x, \lambda)$ $>0$ for $x \in \Sigma(\lambda)$.

In $(-a<\lambda<a)$ there is a maximal open interval $(-a<\lambda<\mu)$ for which the inequality

$$
\begin{equation*}
w(x, \lambda) \geqslant 0 \quad \forall x \in \Sigma(\lambda) \tag{4.17}
\end{equation*}
$$

holds. We wish to show that $\mu=a$. Suppose $\mu<a$ - we will obtain a contradiction. By continuity we have $w(x, \mu) \geqslant 0$ for $x \in \Sigma(\mu)$, and as before we infer that

$$
\begin{equation*}
w(x, \mu)>0 \quad \text { for } \quad x \in \Sigma(\mu) \tag{4.17}
\end{equation*}
$$

By definition of $\mu$ there is a sequence $\lambda^{i} \forall \mu$ and points $x^{i} \in \Sigma\left(\lambda^{i}\right)$ such that

$$
\begin{equation*}
\underline{u}\left(x^{i}\right)>u\left(x_{1}^{i}+a-\lambda^{i}, y^{i}\right)=u\left(\tilde{x}^{i}\right) . \tag{4.18}
\end{equation*}
$$

Here we have set $x^{i}+\left(a-\lambda^{i}\right) e_{1}=\tilde{x}^{i}$. We may suppose that in $\overline{\Sigma\left(\lambda^{i}\right)}$, $w$ has its (negative) minimum at $x^{i}$. So $\nabla_{x} w=0,\left\{w_{j k}\right\} \geqslant 0$ there. A suitable subsequence $x^{i}$ converges to a point $\bar{x}$ in $\overline{\Sigma(\mu)}$, with $\widetilde{x}^{i} \rightarrow \bar{x}+(a-\mu) e_{1}=\bar{x}$. Hence

$$
\begin{equation*}
w(\bar{x}, \mu)=0 \tag{4.19}
\end{equation*}
$$

and so $\bar{x} \in \partial \Sigma(\mu)$. At $(\bar{x}, \mu), \nabla_{x} w=0,\left\{w_{j k}\right\} \geqslant 0$.
In $(x, \lambda)$ space we set

$$
V=U \cap\{\lambda<\mu\}
$$

Several cases can occur and each has to be treated. In the following we usually write $w(x)$ to represent $w(x, \mu)$.

Case 1. $-a<\bar{x}_{1}<\mu$. We may suppose $e_{2}$ is exterior normal to $\partial \Omega$ at $\bar{x}$. In $V$ we may apply Lemma $\tilde{H}$ (i.e. Lemma 4.1) and conclude that

$$
\begin{equation*}
w_{2}(\bar{x}, \mu)<0, \text { contradiction } \tag{4.20}
\end{equation*}
$$

Case 2. $-a=\bar{x}_{1}, \bar{y} \in \omega$. Applying Lemma $\tilde{\mathrm{H}}$ again we see that $w_{1}(\bar{x})>0-$ again a contradiction.

Case 3. $\bar{x}=(\mu, \bar{y}), \bar{y} \in \omega$. By Lemma $\tilde{\mathrm{H}}, w_{1}(\bar{x})<0-$ contradiction.
Case 4. $\bar{x}=(-a, \bar{y}), \bar{y} \in \partial \omega$. This is treated as in the proof of Theorem 2.1. We know

$$
\begin{equation*}
\nabla_{x} w=0,\left\{w_{j k}\right\} \geqslant 0 \text { at } \bar{x} \tag{4.21}
\end{equation*}
$$

Furthermore, since in $\bar{V}$, $w$ has a minimum, zero, at $(\bar{x}, \mu)$, we have there $\partial_{\lambda} w \leqslant 0$.
It follows from (4.14) that $\left\{w_{j k}\right\}=0$ there.
In $V$ we now apply Lemma $\widetilde{\mathbb{S}}$ (Lemma 4.3). Here near $(\bar{x}, \mu)$ the function $\rho=\rho\left(x_{2}, \ldots, x_{n}\right), \rho<0$ describes $\omega$, and $\sigma=-x_{1}-a$. So $a_{i j} \rho_{i} \sigma_{j}=0$ everywhere. We may suppose ( $1,0, \ldots, 0$ ) is exterior normal to $\partial \omega$ at $\bar{y}$. By the lemma,

$$
\left(\partial_{1}-\partial_{2}\right)^{2} w>0 \text { at } \bar{x}-\text { contradiction }
$$

Case 5. $\bar{x}=(\mu, \bar{y}), \bar{y} \in \partial \omega$. Here $\bar{x}=(a, \bar{y})$. Just proceed as in Case 4 using Lemma $\widetilde{S}$ in $V$ at $(\bar{x}, \mu)$.

Thus all cases are impossible and we conclude that $\mu=a$ - so the Theorem is proved.

The proof of Theorem 4.1 yields also the following.
THEOREM 4.1'. Let $f$ be as in Theorem 4.1. Let $\underline{u}, u \in C^{2}(\bar{\Omega})$ be solutions of (4.1) satisfying

$$
\begin{aligned}
& \underline{u}\left(x_{1}, y\right) \leqslant u\left(x_{1}^{\prime}, y\right) \text { for } x_{1} \leqslant x_{1}^{\prime}, y \in \partial \omega \\
& \underline{u}\left(x_{1}, y\right) \leqslant u(a, y), \underline{u}(-a, y) \leqslant u\left(x_{1}, y\right) \text { for }-a<x_{1}<a, y \in \omega \\
& \text { and } \forall x_{1} \text { in }(-a, a), \exists y \in \omega \text { such that } \\
& \underline{u}(-a, y)<u\left(x_{1}, y\right)
\end{aligned}
$$

Then $\underline{u} \leqslant u$ in $\Omega$.

A special case of this is

THEOREM 4.1". Let $u$ and $\underline{u}$ be positive functions on $\Omega=(-a, a)$ belonging to $C^{2}(\Omega) \cap C(\bar{\Omega})$ and both satisfying

$$
\begin{aligned}
& \ddot{u}+f(x, u)=0 \text { in } \Omega \\
& u=0 \text { at } \pm a .
\end{aligned}
$$

Here $f$ is continuous in $(x, u)$, Lipschitz continuous in $u$, and $f$ is symmetric in $x$, and nondecreasing in $x$ for $-a<x<0$. Then
(i) $u$ and $\underline{u}$ are symmetric in $x$ and $u_{x}, \underline{u}_{x}>0$ on $-a<x<0$.
(ii) The functions $u, \underline{u}$ are identical or one is strictly greater than the other in $\Omega$.

Proof: (i) is proved in [4]; (ii) follows from Theorem 4.1'.

This result does not hold in higher dimensions, for $n \geqslant 3$ see $\mathrm{Lin}, \mathrm{Ni}[10]$.
Let us turn now to noncylindrical domains. Consider a bounded domain $\Omega$ in $R^{n}$ with smooth boundary and which is convex in the $x_{1}$ direction. Can one extend Theorem 4.1 to this? Under somewhat stronger conditions one can give a very simple proof. Here is such a result.

THEOREM 4.2. Let $\Omega$ be as just described, and assume that $u \in C^{2}(\Omega) \cap C(\bar{\Omega})$ satisfies (4.1) and (4.2) with $f$ as in Theorem 4.1. Assume that if $\left(x_{1}^{\prime}, y\right)$, $\left(x_{1}^{\prime \prime}, y\right) \in \partial \Omega, \quad x_{1}^{\prime}<x_{1}^{\prime \prime}$, then

$$
\begin{equation*}
\phi\left(x_{1}^{\prime}, y\right)<\phi\left(x_{1}^{\prime \prime}, y\right) \tag{4.22}
\end{equation*}
$$

and if in addition $(x, y) \in \Omega, x_{1}^{\prime}<x_{1}<x_{1}^{\prime \prime}$, then

$$
\begin{equation*}
\phi\left(x_{1}^{\prime}, y\right)<u(x, y)<\phi\left(x_{1}^{\prime \prime}, y\right) . \tag{4.23}
\end{equation*}
$$

Then the conclusion of Theorem 4.1 holds.
Proof: We may suppose that the longest open interval in $\Omega$ parallel to $e_{1}$ has length 2 a and $x_{1}= \pm a$ at its right and left endpoints. For $-a<\lambda<a$ let

$$
\Sigma(\lambda)=\left\{x \in \Omega ; x+(a-\lambda) e_{1} \in \Omega\right\} .
$$

Suppose $\underline{u}$ is a solution of (4.1) in $\Omega$ satisfying the same conditions as $u$. As before it suffices to show (4.6) for $x \in \Sigma(\lambda)$. This then proves the theorem. In the region in $(x, \lambda)$ space:

$$
U=\{(x, \lambda) ; x \in \Sigma(\lambda),-a<\lambda<a\}
$$

$w$ satisfies the parabolic inequality (4.14). On the spatial boundary of $U$, i.e. the part of the boundary in $\lambda<a$, we have $w \geqslant 0$. In fact for $-a<\lambda<a$ by (4.22),

$$
\begin{equation*}
w(x, \lambda)>0 \text { if } x \in \partial \Sigma(\lambda), x \text { and } x-(a-\lambda) e_{1} \in \partial \Omega . \tag{4.24}
\end{equation*}
$$

In $x$-space, on the boundary of each components of $\Sigma(\lambda)$ there is a point $x$ with $x$ and $x-(a-\lambda) e_{1} \in \partial \Omega$, otherwise $\partial \Sigma(\lambda)$ would consist entirely of $\partial \Omega$ - impossible. So

$$
\begin{equation*}
w \geqq 0 \text { for } x \in \partial \Sigma(\lambda) \text {. } \tag{4.25}
\end{equation*}
$$

For $0<a+\lambda$ small, Proposition 1.1 implies that

$$
w(x, \lambda) \geqslant 0 .
$$

Then by (4.25) and Lemma 4.2, it follows that $w(x, \lambda)>0$ in $U$ for $0<a+\lambda$ small.

In $(-a<\lambda<a)$ there is a maximal open interval $(-a<\lambda<\mu)$ for which the inequality (4.17) holds, and we wish to show that $\mu=a$. Suppose $\mu<a$. By continuity $w(x, \mu) \geqslant 0$ for $x \in \Sigma(\mu)$, and again we conclude that

$$
w(x, \mu)>0 \text { for } x \in \Sigma(\mu)
$$

Proceed as before: by definition of $\mu$ there is a sequence $\lambda^{i} \downarrow \mu$ and points $x^{i} \in \Sigma\left(\lambda^{i}\right)$ such that (4.18) holds. And we may suppose that in $\overline{\Sigma\left(\lambda^{i}\right)}, w$ has its (negative) minimum at $x^{i}$. So $\nabla_{x} w=0,\left\{w_{j k}\right\} \geqslant 0$ there. A subsequence $x^{i} \rightarrow$ $\rightarrow \bar{x} \in \overline{\Sigma(\mu)}$, because of (4.22), (4.23) $\tilde{x}^{i} \rightarrow \bar{x}=\bar{x}+(a-\mu) e_{1}$. Clearly $w=0$, $\nabla_{x} w=0,\left\{w_{j k}\right\} \geqslant 0$ at $(\bar{x}, \mu)$. So $\bar{x} \in \partial \Sigma(\mu)$.

In virtue of (4.23) and (4.24),

$$
w>0 \text { on } \partial \Sigma(\mu), \text { contradiction. }
$$

So in this case we know $\mu=a$ and Remark 4.2 is proved.

As before we see that Corollary 4.1 holds. In addition we have the analogue of Cor. 4.2: antisymmetry.

It is natural to ask if one can relax condition (4.22). Here is a result in this direction.

THEOREM 4.3. Let $\Omega$ be as in Theorem 4.2. In place of (4.22), assume for $\left(x_{1}^{\prime}, y\right)$, $\left(x_{1}^{\prime \prime}, y\right) \in \partial \Omega, x_{1}^{\prime}<x_{1}^{\prime \prime}$

$$
\begin{equation*}
\phi\left(x_{1}^{\prime}, y\right) \leqslant \phi\left(x_{1}^{\prime \prime}, y\right) \tag{4.22}
\end{equation*}
$$

and if in addition, $(x, y) \in \Omega, x_{1}^{\prime}<x_{1}<x_{1}^{\prime \prime}$, then strict inequality holds in (4.22)', and (4.23) holds. Assume in addition Hypothesis 3): if the segment joining $\left(x_{1}^{\prime}, y\right)$ to $\left(x_{1}^{\prime \prime}, y\right)$ lies on $\partial \Omega$ ard on it $\phi$ is constant, then the exterior normal $\nu$ to $\partial \Omega$ is also constant on it. (This automatically holds if $n=2$ ). The conclusion of Theorem 4.1 then holds.

Proof: Proceed as before. Suppose $\mu<a$. Again we obtain a point $\bar{x} \in \partial \Sigma(\mu)$ such that at $(\bar{x}, \mu), w=0, \nabla_{x} w=0,\left\{w_{j k}\right\} \geqslant 0$. Furthermore $\partial_{\lambda} w \leqslant 0$ at this point so by (4.14) $\left\{w_{j k}\right\}=0$ there. $w(x, \mu)>0$ for $x \in \Sigma(\mu)$. If $\bar{x}$ belongs to a smooth part of $\partial \Sigma(\mu)$ we have, as before, using Lemma $\widetilde{\mathrm{H}}, \partial_{\nu} w(\bar{x}, \mu)<0-$ contradiction. So $\bar{x}$ is such that $\bar{x}$ and $\bar{x}+(a-\lambda) e_{1} \in \partial \Omega$ and $\phi$ has the same value at these points. By conditions (4.22)' and (4.23) it follows that the segment joining these points belongs to $\partial \Omega$.

But then by Hypothesis 3), the exterior normal $\nu$ to $\partial \Omega$ is constant on this segment; we may suppose $\nu=e_{2}$. Because of this, we may apply Lemma 4.4 and get $\partial_{\nu} w(\bar{x}, \mu)<0$ - contradiction. Thus $\mu=a$ and the theorem is proved.

Suppose the domain $\Omega$ is symmetric in $x_{1}$ about $x_{1}=0$, and of course, convex in the $x_{1}$ direction, and suppose $\phi$ satisfies the conditions in Theorem 4.3. Then Hypothesis 3 ) is automatically satisfied. Thus we have

COR. 4.4. ANTISYMMETRY. Assume the conditions of Theorem 4.3, without, however, Hypothesis 3). Suppose $\Omega$ is symmetric in $x_{1}$ and that $\phi$ is odd in $x_{1}$. Assume $f(x, u, p)$ is odd in $\left(x_{1}, u, p_{2}, \ldots, p_{n}\right)$. Then $u$ is antisymmetric in $x_{1}$.

## Appendix. Proofs of Proposition 1.1, of Lemmas 4.1-4.4 et al.

Proposition 1.1 and Lemmas 4.1, 4.2 are well known in case $\beta>0$ (see e.g. Chapter 3, Sec. 2 of Protter-Weinberger [7]). The proofs for our case $\beta \geqslant 0$ are very similar and will just be sketched. In particular the proof of Proposition 1.1 is essentially the same as the corollary on page 213 of [4].

Proof of Proposition 1.1.: Suppose $b_{1} \geqslant-b, c \leqslant c_{1}$ with $b, c_{1} \geqslant 0$, and suppose $\epsilon>0$ so small that

$$
c_{1} \exp \left(4 b \epsilon / c_{0}\right)<c_{1}+2 b^{2} / c_{0} .
$$

Recall that $a$ is the ellipticity constant in (1.8). The function

$$
g=e^{\alpha(-a+2 \epsilon)}-e^{\alpha x_{1}}
$$

is positive in $\bar{V}$ and satisfies

$$
-L g=\left(a_{11} \alpha^{2}+b_{1} \alpha\right) e^{\alpha x_{1}}-c\left(e^{\alpha(-a+2 \epsilon)}-e^{\alpha x_{1}}\right)
$$

Thus choosing $\alpha=2 b / c_{0}$ we see that

$$
\begin{aligned}
& -e^{-\alpha x_{1}} L g \geqslant c_{0} \alpha^{2}-b \alpha+c_{1}-c_{1} e^{2 \alpha \epsilon} \\
& =\frac{2 b^{2}}{c_{0}}+c_{1}-c_{1} e^{2 \alpha \epsilon}, \text { since } \alpha=2 b / c_{0} \\
& >0
\end{aligned}
$$

So the function

$$
v=w / g
$$

satisfies

$$
\left(L^{\prime}+\frac{L g}{g}-\beta \partial_{t}\right) v>0 \text { in } V
$$

where $L^{\prime}$ is an elliptic operator with no zero order term. Because $L g / g<0$ in $\bar{V}$ we see that $v$ can attain a positive maximum only on $J=\overline{\partial V \cap\{t<T\}}$. So $v \leqslant 0$ and hence $w \leqslant 0$.

Note that for all $\left(x^{0}, T\right)$ belonging to $\partial V \backslash J$ there is an open half ball centred at $\left(x^{0}, T\right)$ lying in $V$.

Proof of Lemma 4.1: By considering a smaller parabolic cap touching $\bar{Q}$ at $(\bar{x}, \bar{T})$ we may suppose that $w<0$ at every point of $\bar{Q}$ except $(x, T)$. We may suppose $x^{0}=0$ and $\bar{x}=\left(\bar{x}_{1}, 0, \ldots, 0\right), \bar{x}_{1}>0$. In $\bar{Q}$, near $(\bar{x}, T)$, we will construct a $C^{2}$ function $h$, positive in $Q, h=0$ on the curved part of $\partial Q, h_{1}<0$ at $(\bar{x}, \bar{T})$ and satisfying

$$
\left(L-\beta \partial_{t}\right) h>0
$$

in the region $\tilde{Q}=Q \cap\left\{x_{1}>\bar{x}_{1} / 2\right\}$. It follows that for $\epsilon>0$ sufficiently small,

$$
\left(L-\beta \partial_{\tau}\right)(w+\epsilon h)>0
$$

in $\widetilde{Q}$, and on the boundary of $\widetilde{Q}$ we have $w+\epsilon h \leqslant 0$, except the top, $t=\bar{T}$.
Now the cap $Q$ may be chosen as small as we like, so we find from Proposition 1.1 that $w+\epsilon h \leqslant 0$ in $\tilde{Q}$. Since $w+\epsilon h=0$ at $(\bar{x}, T)$, necessarily $\partial_{\nu}(w+\epsilon h) \geqslant 0$ there. But $\partial_{\nu} h<0$ there - the result follows.

Now to construct $h$. Set

$$
h=e^{-\alpha\left(|x|^{2}-t\right)}-e^{-\alpha(\delta-T)}
$$

For $\alpha$ positive, this is positive in $P$ and zero on $\partial P$, and $\nabla_{x} h \neq 0$ on $\partial P$ except at the bottom point ( $x^{0}, T-\delta$ ). Furthermore

$$
\begin{aligned}
& \left(L-\beta \partial_{t}\right) h=e^{\alpha\left(t-|x|^{2}\right)}\left[4 \alpha^{2} a_{i j} x_{i} x_{j}-2 \alpha \Sigma a_{i i}-2 \alpha b_{i} x_{i}+c-\beta \alpha\right]- \\
& -c e^{-\alpha(\delta-T)}>0 \text { in } \tilde{Q} \text { for } \alpha \text { large }
\end{aligned}
$$

The lemma is proved.
Proof of Lemma 4.2.: We may suppose $t^{0}=T$ : It suffices to show that $w<0$ on every polygonal path in $V_{T}$ starting at $\left(x^{0}, T\right)$. We show that $w<0$ on the first closed straight segment of this path starting at $\left(x^{0}, T\right)$. Continuing this argument one finds that $w<0$ on the whole path.

We may suppose $x^{0}=0$ and the segment is $S=\left(x_{1}, 0, \ldots, 0, T\right), 0 \leqslant x_{1} \leqslant \ell$. Consider a small quarter ball lying in $V$ in the closure of which $w<0$ :

$$
\left\{(x, t) ; x_{1}^{2}+\ldots+x_{n}^{2}+(T-t)^{2}<\epsilon^{2}, t \leqslant T, x_{1}>0\right\}
$$

Now stretch this in the $x_{1}$ direction, as a quarter ellipsoid:

$$
E_{\tau}=\left\{(x, t) ; t \leqslant T, x_{1}>0, \tau x_{1}^{2}+x_{2}^{2}+\ldots+x_{n}^{2}+(T-t)^{2}<\epsilon\right\}
$$

by decreasing $\tau$ from 1. As we keep decreasing $\tau$ to some $\tau_{0}, E_{\tau_{0}}$ will eventually contain the whole segment $S$ and will still belong to $\tilde{V}$ if $\epsilon$ is sufficiently small. We claim that for all $\tau, 1 \geqslant \tau \geqslant \tau_{0}$ we have $w<0$ on $\bar{E}_{\tau}$ - and so $w<0$ on $S$. If not there is a first such $\tau$ such that $w<0$ in $E_{\tau}$ and $w=0$ at some point $z$ on the curved boundary of $E_{\tau}$. This point is not the bottom point ( 0 , $T-\epsilon$ ) and so we can fit a small parabolic cap in $E_{\tau}$ touching $E_{\tau}$ at the point z. But by Lemma 4.1 $\nabla w \neq 0$ at $z$. Impossible since $z$ is a point in $\tilde{V}$.

Proof of Lemma 4.3. We will follow the proof of Lemma S of [4] - the reader should have that paper in hand. We may assume $x^{0}=0, T=0$.

The conditions of the lemma are invariant under changes of variable $(x, t) \rightarrow$ $\rightarrow(y(x, t), t)$. If $\rho_{1}, \sigma_{n} \neq 0$ at $(0,0)$ we may introduce new spatial variables ( $\left.\rho, x_{2}-\chi_{2}(t), \ldots, x_{n-1}(t)-\chi_{n-1}(t), \sigma\right)$. The transformed $V$, near the origin, consists of the region, $x_{1}, x_{n}, t<0$. The curve $\mathscr{C}$ now lies on the $t$ axis. Next, by considering a slightly smaller region $x_{1}+\gamma \Sigma_{2}^{n-1} x_{\alpha}^{2}, x_{n}+\gamma \Sigma_{2}^{n-1} x_{\alpha}^{2}, t<0$, here $\gamma>0$, we have $w<0$ on its closure, near ( 0,0 ), except possibly on the $t$ axis. Making a second change of variables $(x, t) \rightarrow(y, t), y_{j}=x_{j}+\gamma \Sigma_{2}^{n-1} x_{\alpha}^{2}$, $j=1$ and $n, y_{\alpha}=x_{\alpha}, 2 \leqslant \alpha \leqslant n-1$, we obtain: $V=\left\{x_{1}, x_{n}, t<0\right\}$ near the origin, and $w<0$ in $\bar{V}$ except possibly on the negative $t$ - axis. $\mathscr{C}$ is still a segment on the nonpositive $t$-axis. We're still not through with changes of variables. We wish to have

$$
\begin{equation*}
a_{1 \alpha}=a_{\alpha n}=0 \text { on } \mathscr{C} \text { for } 1<\alpha<n \tag{A.1}
\end{equation*}
$$

As in [4] this is achieved by a change of variables: $y=x_{1}, y_{n}=x_{n}, y_{\alpha}=x_{\alpha}+$ $+c_{\alpha}(t) x_{1}+d_{\alpha}(t) x_{n}$; see the computation there.

So finally we have near the origin: $V=\left\{x_{1}, x_{n}, t<0\right\}$, $\mathscr{C}$ is the nonpositive $t$ - axis, and (A.1) holds.

Now we follow pages 240-242 of [4]. There one took $c=0$. We will not do that here; we write

$$
L=M+c
$$

The inequalities expressed in pages 240-242 of [4] will hold for $M$. There, two functions

$$
\phi=x_{1}+k \sum_{2}^{n-1} x_{\beta}^{2}, \psi=x_{n}+k \sum_{2}^{n-1} x_{\beta}^{2}
$$

were introduced, and a region, in $x$-space, near the origin,

$$
G=\{\phi, \psi<0\}
$$

Now we consider the region in $(x, t)$ space near the origin:
(A.2)

$$
\widetilde{G}=\{\phi, \psi, t<0\}
$$

On pages 240,24] of [4] for large $k$ and then large $\alpha$, the functions

$$
z(x)=g h, g=e^{-\alpha \phi}-1, h=e^{-\alpha \psi}-1
$$

were constructed. In a small region $G_{\alpha, \delta}=\{-1 / \alpha \leqslant \phi, \psi<0\} \cap\{|x|<\delta\}$ the function $z$ was shown to satisfy

$$
\begin{equation*}
z>0, \text { and } z=0 \text { on } \phi=0 \text { and on } \psi=0 \tag{A.3}
\end{equation*}
$$

$$
\begin{equation*}
\partial_{\nu} z=0, \partial_{\nu}^{2} z>0 \text { at } 0 \tag{A.4}
\end{equation*}
$$

and
(A.5) $\quad \frac{e^{\alpha(\phi+\psi)}}{\alpha^{2}} M z \geqslant c_{1} \alpha(|\phi|+|\psi|), c_{1}>0$.

In verifying (A.4) the properties

$$
\begin{equation*}
a_{1 n} \geqslant C\left(x_{1}+x_{n}-\sum_{2}^{n-1} x_{\beta}^{2}\right), C>0 \tag{A.6}
\end{equation*}
$$

$$
\begin{equation*}
\left|a_{1 \beta}\right|,\left|a_{n \beta}\right| \leqslant C|x|, 1<\beta<n, \tag{A.7}
\end{equation*}
$$

were used. Because of (4.10) and (A.1), these properties (A.6), (A.7) continue to hold in the region

$$
\widetilde{G}_{\alpha, \delta}=\left\{(x, t) ; x \in G_{\alpha, \delta},-\delta<t<0\right\}
$$

for $\delta$ sufficiently small. Thus (A.3) and (A.5) hold in $\widetilde{G}_{\alpha, \delta}$.
Consider now in $\widetilde{G}_{\alpha, \delta}$ the function

$$
\tilde{z}=z+t z^{1 / 2}
$$

In the region $W$ where it is positive, we have

$$
\begin{aligned}
& \left(L-\beta \partial_{t}\right) \widetilde{z}=M z+c z+\frac{t}{2} z^{-1 / 2} M z-\frac{t}{4} z^{-3 / 2} a_{i j} z_{i} z_{j} \\
& +c t z^{1 / 2}-\beta z^{1 / 2} \\
& \geqslant \frac{1}{2} M z+c z-C z^{1 / 2}
\end{aligned}
$$

since $t<0$, hence

$$
\geqslant \frac{c_{1}}{2} \alpha^{3}(|\phi|+|\psi|)-c\left(|z|+|z|^{1 / 2}\right) \quad \text { by } \quad \text { (A.5). }
$$

Now in $\tilde{G}_{\alpha, \delta}$ we have

$$
\begin{aligned}
& z=\left(e^{-\alpha \phi}-1\right)\left(e^{-\alpha \psi}-1\right) \\
& \leqslant e^{2} \alpha^{2} \phi \psi
\end{aligned}
$$

since $e^{s}-1 \leqslant e s$ for $0<s \leqslant 1$. Hence, with a different constant $C$,

$$
\begin{aligned}
& \left(L-\beta \partial_{t}\right) \tilde{z} \geqslant \frac{c_{1}}{2} \alpha^{3}(|\phi|+|\psi|)-C \alpha^{2} \phi \psi-C \alpha(\phi \psi)^{1 / 2} \\
& >0 \text { in } W
\end{aligned}
$$

for $\alpha$ large.
Now we use the function $\tilde{z}$ as a comparison function in the usual way. On $\partial W$ except at the origin we have $w<0$. Hence for small $\epsilon>0$ we have

$$
w+\epsilon \tilde{z} \leqslant 0
$$

on the entire boundary of $W$ lying in $t<T$. For $\alpha$ large, $W$ is as narrow as we like in the $x_{1}$ direction. Thus Proposition 1.1 applies and we conclude that

$$
w+\epsilon \tilde{z} \leqslant 0 \text { in } W .
$$

But $w+\epsilon \tilde{z}=0$ at $(0,0)$, so $\partial_{\nu}(w-\epsilon \widetilde{z}) \geqslant 0$ or $\partial_{\nu}^{2}(w+\epsilon \widetilde{z}) \leqslant 0$. Since $\tilde{z}$ satisfies (A.4):

$$
\partial_{\nu} \tilde{z}=0, \quad \partial_{\nu}^{2} \tilde{z}>0
$$

the desired result (4.11) follows.

Lemma 4.4 is a special case of the following parabolic analogue of Lemma 1.1 of [4] (with $\mu>0$, so $\theta_{0}>\pi / 2$, and $p<2$ ). We take $(\bar{x}, T)=(0,0)$.

LEMMA A.1. Let $w$ and $V$ be as in Lemma 4.4 with $a_{i j} \in C(\bar{V})$. In place of (4.9)' assume that at $(0,0)$

$$
\Sigma a_{i j} \rho_{i} \sigma_{j}=\mu \sqrt{\Sigma a_{i j} \rho_{i} \rho_{j}} \sqrt{\Sigma a_{i j} \sigma_{i} \sigma_{j}}
$$

for some constant $\mu$; clearly $-1<\mu<1$. Set $\theta_{0}=(\arccos )(-\mu)$. Suppose $p>$ $>\pi / \theta_{0}$. Let C be a closed cone in $\{t=0\}$ with vertex at $(0,0)$ and such that for some $\epsilon>0, \mathrm{C} \cap\{0<|x|<\epsilon\}$ lies in $V_{0}$. Then there is a positive constant $\delta$ and a neighbourhood in C of 0 in which
(A.8)

$$
w+\delta|x|^{p} \leqslant 0
$$

In particular, if $\mu>0$ we may take $p>2$, and it follows that if $w \in C^{1}(\overline{\tilde{V}})$ and $\nabla_{x} w(0,0)=0$ then on any direction $s$ from $(0,0)$ entering $V_{0}$ transversally to the boundary, the second spatial derivatives on $w$ cannot be bounded near the origin. On the other hand if $\mu<0$, and $w \in C^{\left[\pi / \theta_{0}\right], \alpha}$ for $\alpha>\pi / \theta_{0}-\left[\pi / \theta_{0}\right]$, then we may take $p<\left[\pi / \theta_{0}\right]+\alpha$ and conclude that at least one of the derivatives

$$
\left(\frac{\partial}{\partial s}\right)^{j} w, j=1, \ldots,\left[\frac{\pi}{\theta_{0}}\right]
$$

is negative at $(0,0)$.

Proofs: We follow the proof of Lemma A. 1 in [4]. As before we may suppose $w<0$ except possibly on $\{\rho=\sigma=0\}$. Proceeding as in [4] we may arrange that near the origin $V$ has the form

$$
V=\left\{(x, t) ; x_{n}>0, x_{1}>x_{n} \cot \theta_{0}, t<0\right\}
$$

Choose $k=p \theta_{0} / \pi>1$.
Following [4] we also make

$$
a_{11}=a_{n n}=1 \quad \text { at }(0,0)
$$

We have $a_{1 n}=0$ there. With $\zeta=x_{1}+i x_{n}$, in [4] we constructed the functions of $x$ :

$$
\begin{aligned}
& v=\operatorname{Im}\left(\zeta^{\pi / \theta_{0}}\right) \\
& z=v^{k}
\end{aligned}
$$

Near the origin in $V, z$ satisfies

$$
L z \geqslant c_{1}|\zeta|^{p-2}, c_{1}>0,|z|<C|\zeta|^{p}
$$

Near the origin, in the region $W$ where it is positive, consider the function

$$
\widetilde{z}=z+t z^{\ell}
$$

with $1>\ell \geqslant 0, \ell \geqslant 1-2 / p$. (For the case of Lemma 4.4 where $\mu>0$ we may take $p<2$ and $\ell=0$ ). We have

$$
\begin{aligned}
& \left(L-\beta \partial_{t}\right) t z^{\ell}=t \ell z^{\ell-1} L z+t \ell(\ell-1) a_{i j} z_{i} z_{j}+c t z^{\ell}-\beta z^{\ell} \\
& \geqslant t \ell z^{\ell-1} L z-C|\zeta|^{p \ell}
\end{aligned}
$$

since $t(\ell-1) \geqslant 0$. Hence in $W$,

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$$
\begin{aligned}
& \left(L-\beta \partial_{t}\right) \tilde{z} \geqslant\left(1+t \ell z^{\ell-1}\right) L z-C|\zeta|^{p \ell} \\
& \geqslant(1-\ell) L z-C|\zeta|^{p \ell} \\
& \geqslant(1-\ell) c_{1}|\zeta|^{p-2}-C|\zeta|^{p \ell} \\
& >0
\end{aligned}
$$

since $p-2<p \ell$.
For small positive $\epsilon$ we have

$$
L(w+\epsilon \widetilde{z})>0 \text { in } W
$$

and $(w+\epsilon \widetilde{z})<0$ on the boundary of $w$ lying in $t<0$. By Proposition 1.1 it follows ( $W$ is narrow) that

$$
w+\epsilon \tilde{z} \leqslant 0 \text { in } W
$$

On $t=0$ we have therefore

$$
\begin{aligned}
& w \leqslant-\epsilon z \\
& =-\epsilon\left(\operatorname{Im} \zeta^{\pi / \theta_{0}}\right)^{k}
\end{aligned}
$$

Thus on any ray $x=-\tau \nu, \tau>0$, on $t=0$ we have (A.8):

$$
w \leqslant-\epsilon c_{1}|x|^{p}
$$

We conclude with a result used in the proof of Theorem 1.3.
PROPOSITION A.1. Assume Hypotheses 1) and 2) of Section 4. Set

$$
J=\overline{\partial V \cap\{t<T\}}
$$

Suppose that on a set $A$ in $J$, described locally by $\rho=0$ with $\rho \in C^{2}, \nabla_{x} \rho \neq 0$ on $A$, the function $w$ satisfies
(A.9) $\quad \partial_{\nu} w \leqslant 0$
where $\nu(x)$ is a smooth spatial unit vector pointing outside of $V$, with $\nu_{1} \equiv 0$. Assume that $J \backslash A$ is nonempty and that on it, $w \leqslant 0$. If $V$ lies in a sufficiently narrow band $\because 0<a+x_{1}<\epsilon$, then $w \leqslant 0$ in $V$.

Proof: Following the proof of Proposition 1.1 given at the beginning of this section we constructed $g\left(x_{1}\right)>0$ in $\bar{V}$ with $L g<0$. Then

$$
v=w / g
$$

satisfied an inequality of the form (4.7), with new $c<0$ in $\bar{V}$. Since $\nu_{1}=0$ we see that $v$ continues to satisfy $\partial_{v} v \leqslant 0$ on $A$ and $v \leqslant 0$ on $J \backslash A$. Thus it
suffices to prove the proposition in case $c<0$ in $\vec{V}$, which we henceforth assume. We will continue to refer to the solution as $w$ (rather than $v$ ).

Suppose

$$
\max _{\bar{V}} w=K>0
$$

Set $W=w-K$, so $\max W=0 . W$ satisfies

$$
\left(L-\beta \partial_{t}\right) W \geqslant-c K>0 \text { in } V .
$$

It follows that $W$ can achieve its maximum in $\bar{V}$ only on $J$. So $W<0$ in $V$ and $W$ necessarily achieves its maximum at a point on $A$, since on $J \backslash A, W \leqslant-K$. But by Lemma $\widetilde{H}$ (Lemma 4.1) $\partial_{\nu} W>0$ there. Contradiction.

REMARK A.1. It is clear that Proposition A. 1 holds in the elliptic case, i.e. no $t$ present. We may not drop the condition $\nu_{1}=0$.

Consider the simple example in $-\delta<x_{1}<\delta, 0<x_{2}<\pi$ :

$$
w=\cos x_{1} \sin x_{2}
$$

satisfies $\Delta w+2 w=0$, vanishes on $x_{2}=0$ and $x_{2}=\pi$ while on the lateral boundaries $x_{1}= \pm \delta(\delta$ arbitrary small $)$

$$
\partial_{v} w<0
$$

where $\nu$ represents the exterior normal.

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[^0]:    ${ }^{1}$ ) A more general form, for fully nonlinear elliptic equations was given in Theorem 2.1'.

